

Spontaneous creation and dynamics of vortices in Bose-Einstein condensates

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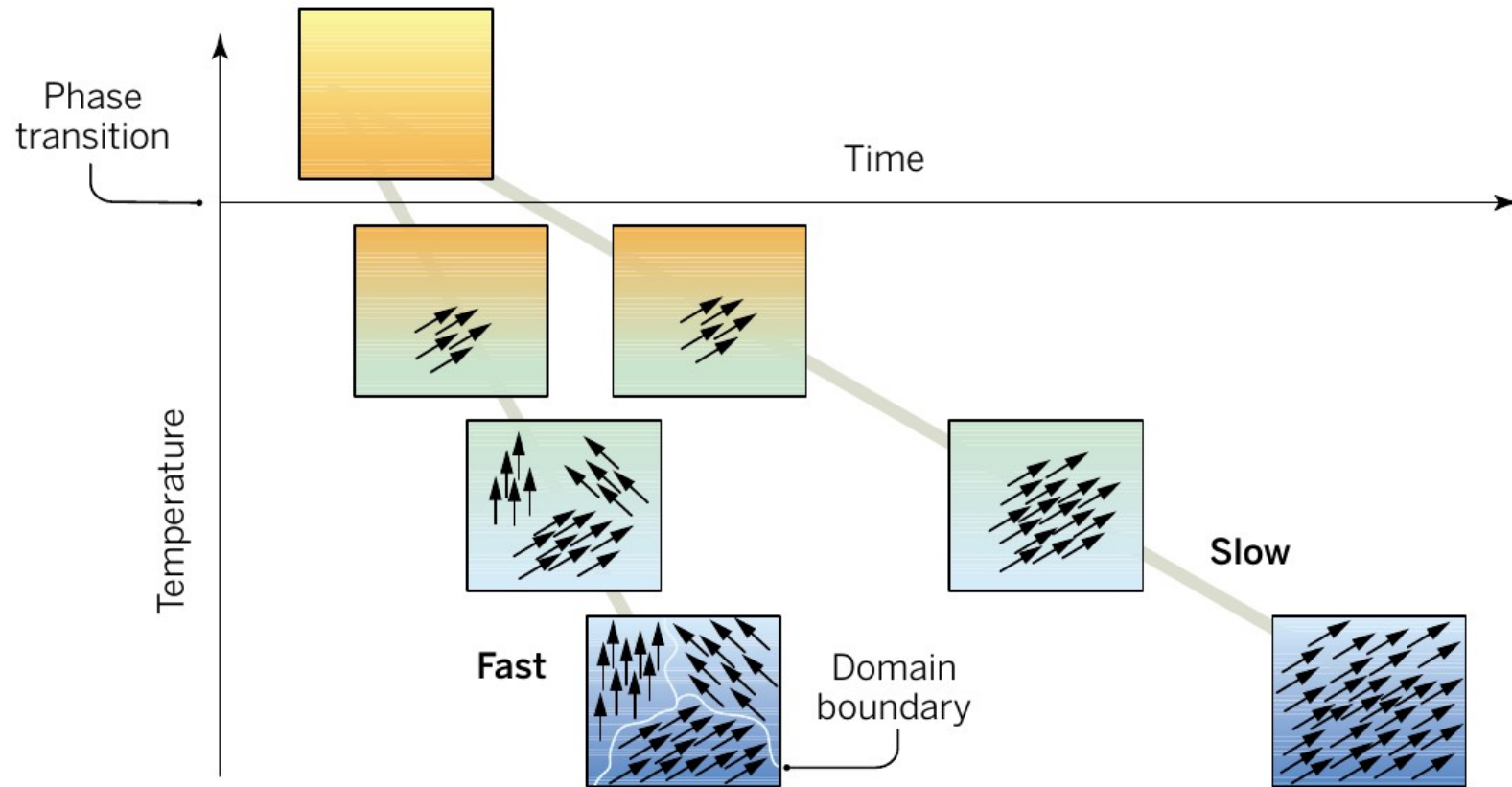
*Journées GdR Atomes Froids et IFRAF
École Normale Supérieure, Paris
5 November 2015*



outline

- Introduction to the Kibble-Zurek mechanism
- Creating defects in Bose condensates via the Kibble-Zurek mechanism
- Defect's characterization
- Dynamics & interactions

the Kibble-Zurek mechanism



- Second-order phase transitions
- Finite rate crossing
- Spontaneous and stochastic production of defects

the Kibble-Zurek mechanism

Power-law scaling

coherence length

$$\xi(t) = \frac{\xi_0}{|\varepsilon(t)|^\nu}$$

relaxation time

$$\tau(t) = \frac{\tau_0}{|\varepsilon(t)|^{z\nu}}$$

reduced parameter:

$$\varepsilon = \frac{\lambda_c - \lambda}{\lambda_c}$$

Case of linear quench

$$\varepsilon(t) = t/\tau_Q$$

time to the transition

$$\tau(t) \approx |\varepsilon/\dot{\varepsilon}|$$

“freezing time”

$$\hat{t} \sim (\tau_0 \tau_Q^{z\nu})^{\frac{1}{1+z\nu}}$$

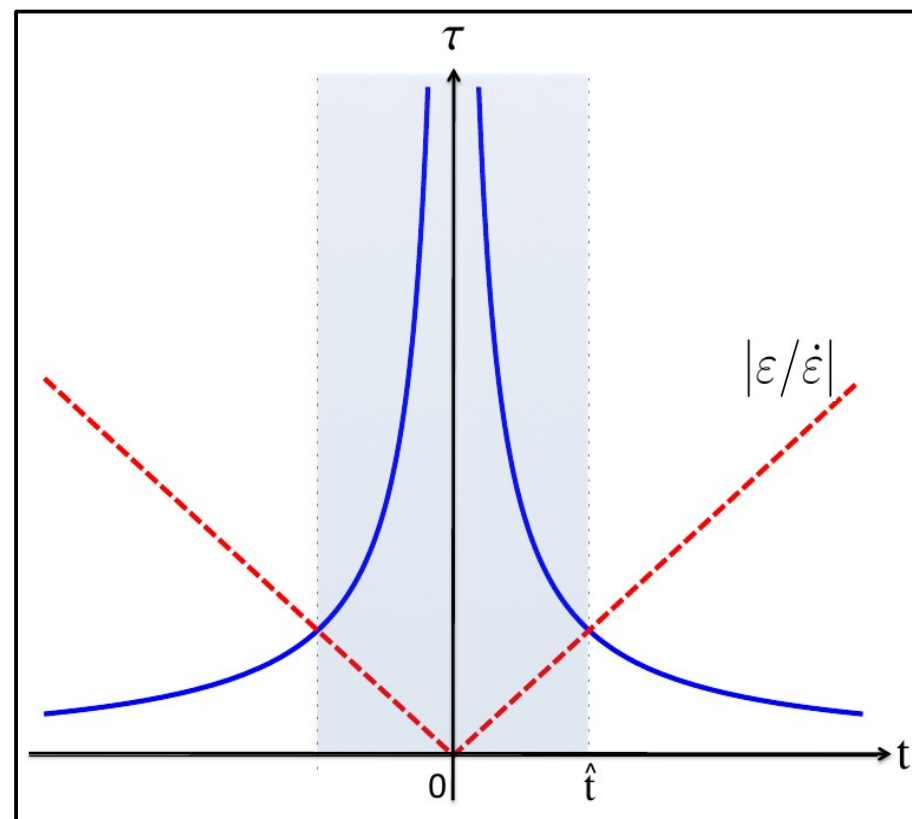
domain size

$$\hat{\xi} = \xi(\hat{t}) = \xi_0 \left(\frac{\tau_Q}{\tau_0} \right)^{\frac{\nu}{1+z\nu}}$$

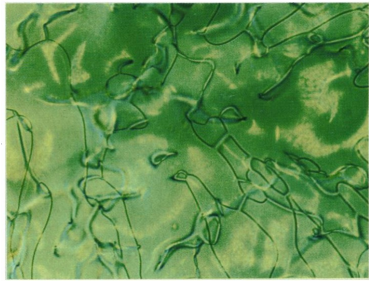
density of defects

$$d \sim \hat{\xi}^{-D}$$

ADIABATIC -- IMPULSE -- ADIABATIC

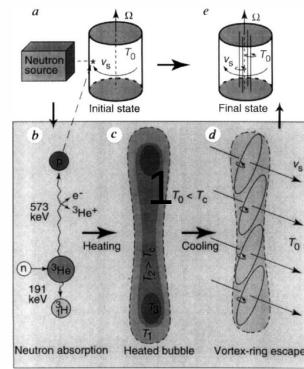


Liquid crystals: isotropic/nematic



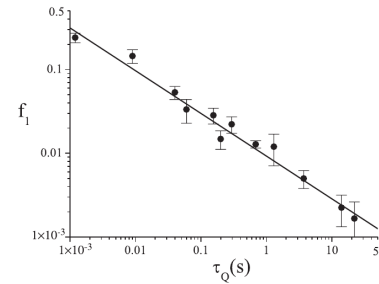
I. Chuang et al. (1991)

Liquid ^3He : normal/SF



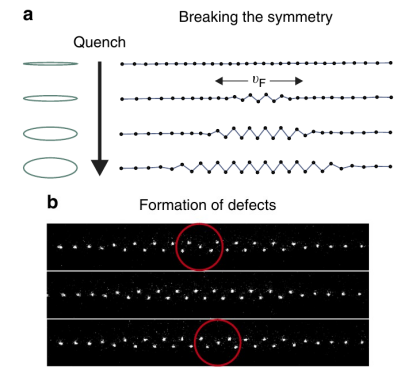
C. Bauerle et al. (1996)
V.M.H Ruutu et al. (1996)

annular Josephson junctions



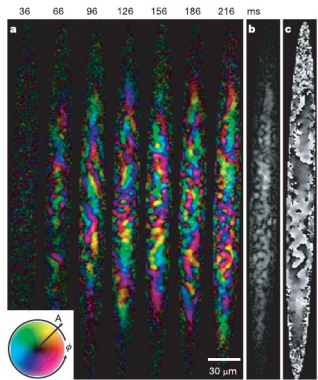
R. Monaco et al. (2009)

1D ion crystals: linear/zig-zag



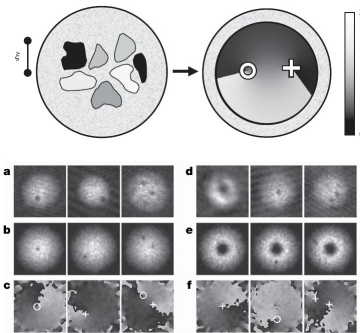
S. Ulm et al. (2013)
K. Pyka et al. (2013)

Bose gases: ferromagnetic



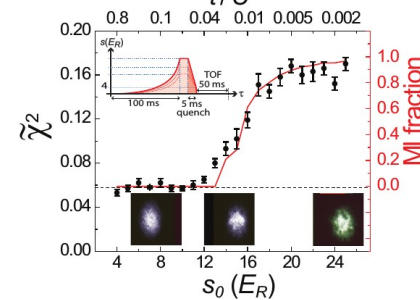
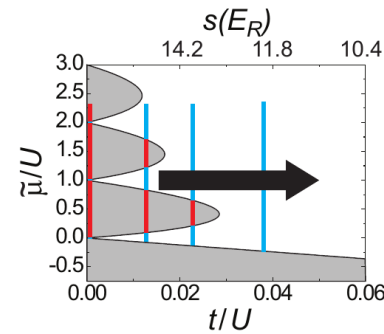
L. E. Sadler et al. (2006)

Bose gases: thermal/BEC



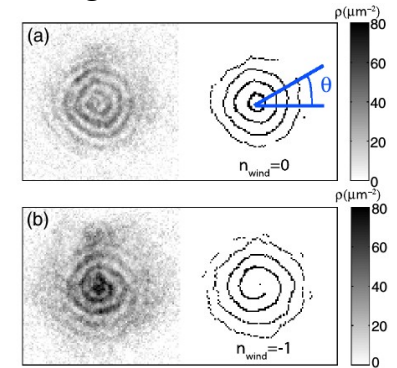
C.N. Weiler et al. (2008)

T=0 Bose gases: Mott/SF



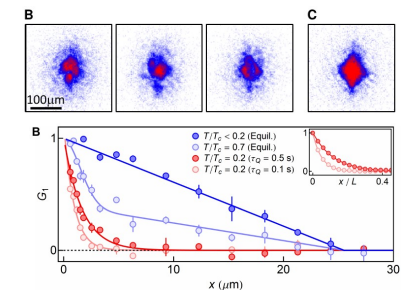
D. Chen et al. (2011)
S. Braun et al. (2014)

Bose gases: thermal/BEC, D < 3



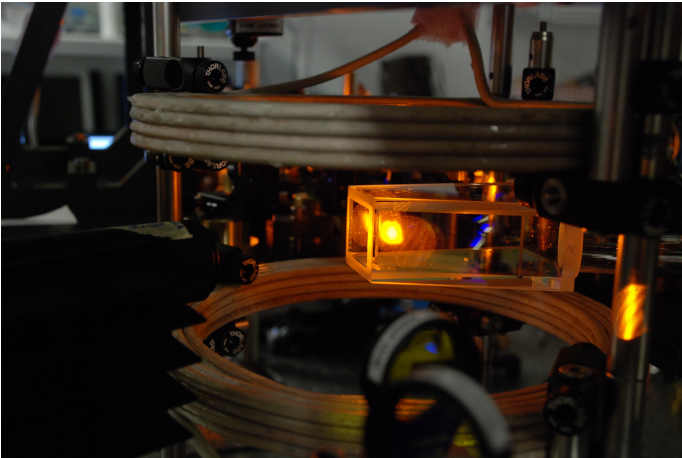
L. Corman et al. (2014)

Hom. Bose gases: thermal/BEC

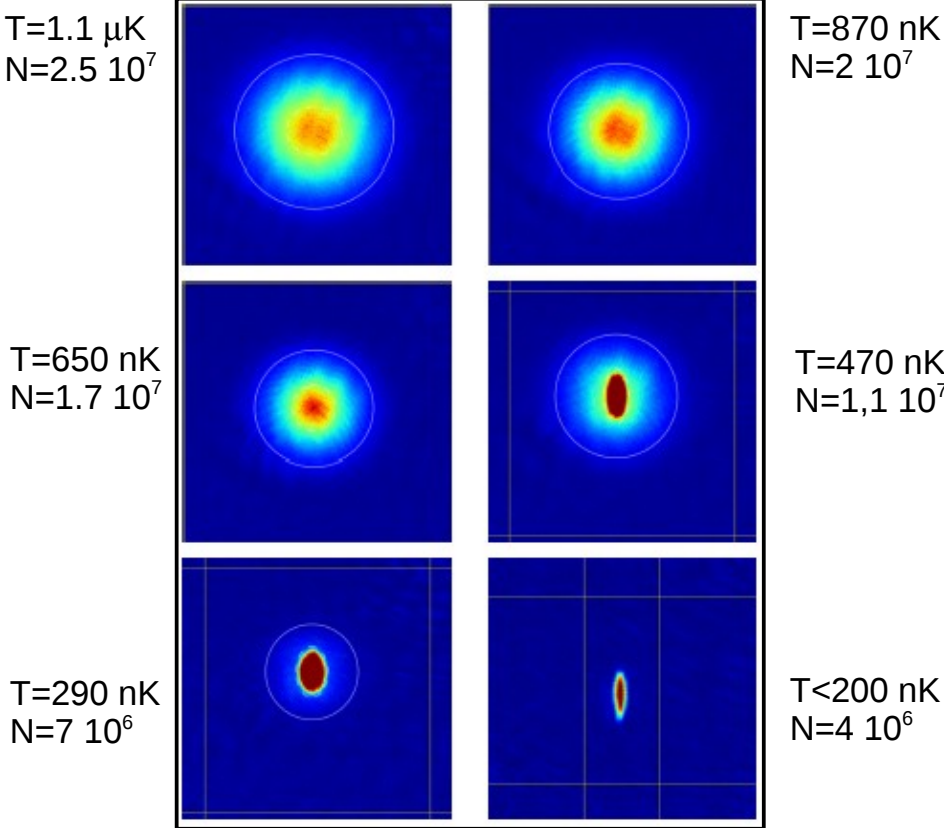
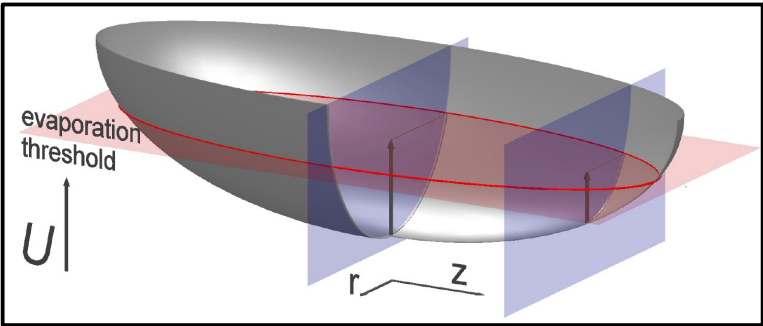


N. Navon et al. (2015)

experimental setup in Trento

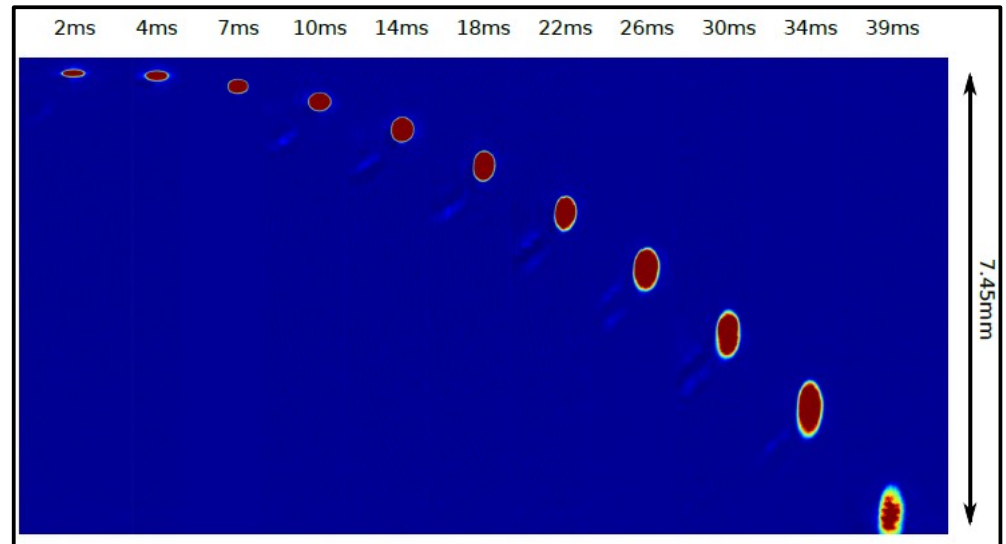


G. Lamporesi *et al.*,
 Rev. Sci. Instrum. **84**, 063102 (2013)

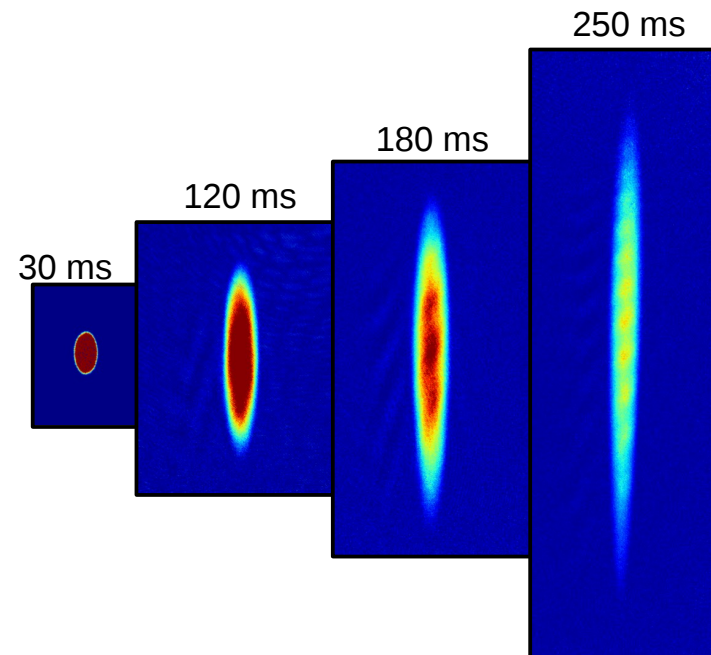


ToF expansion of a BEC

expansion time limited to ~ 40 ms
due to the gravity fall



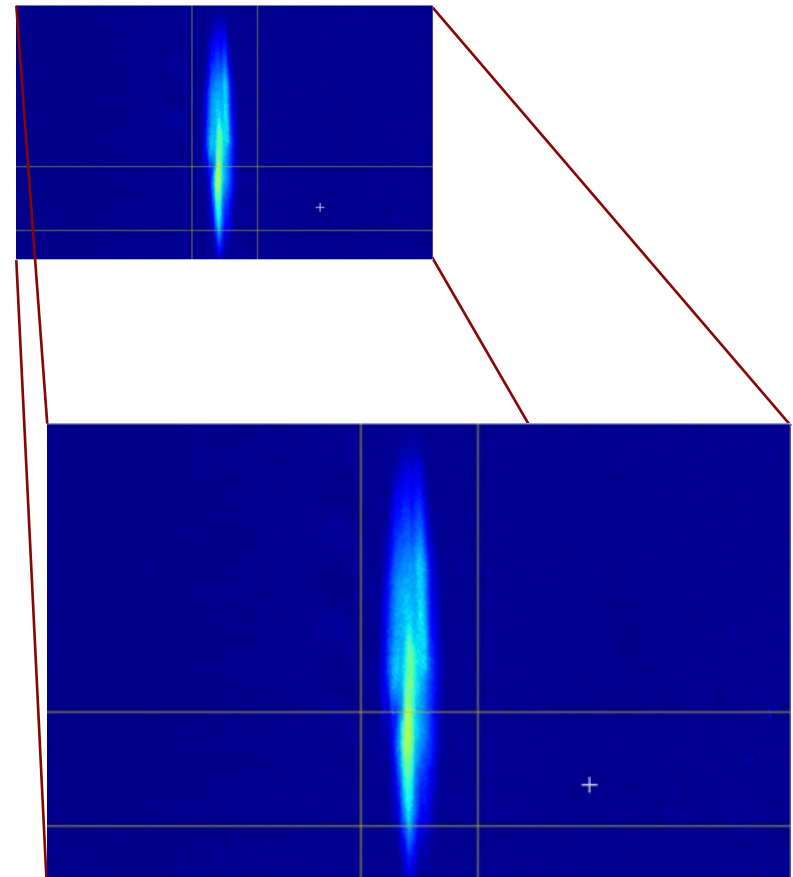
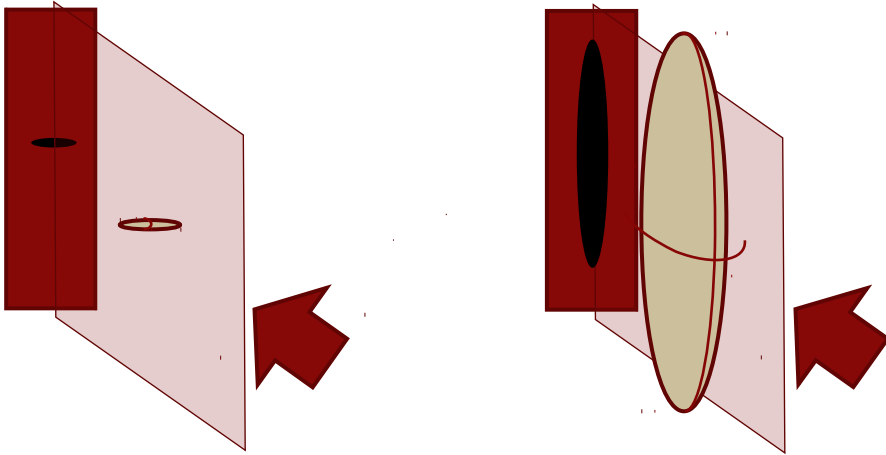
magnetic levitation against the gravity
to increase the expansion time



ToF expansion of a BEC

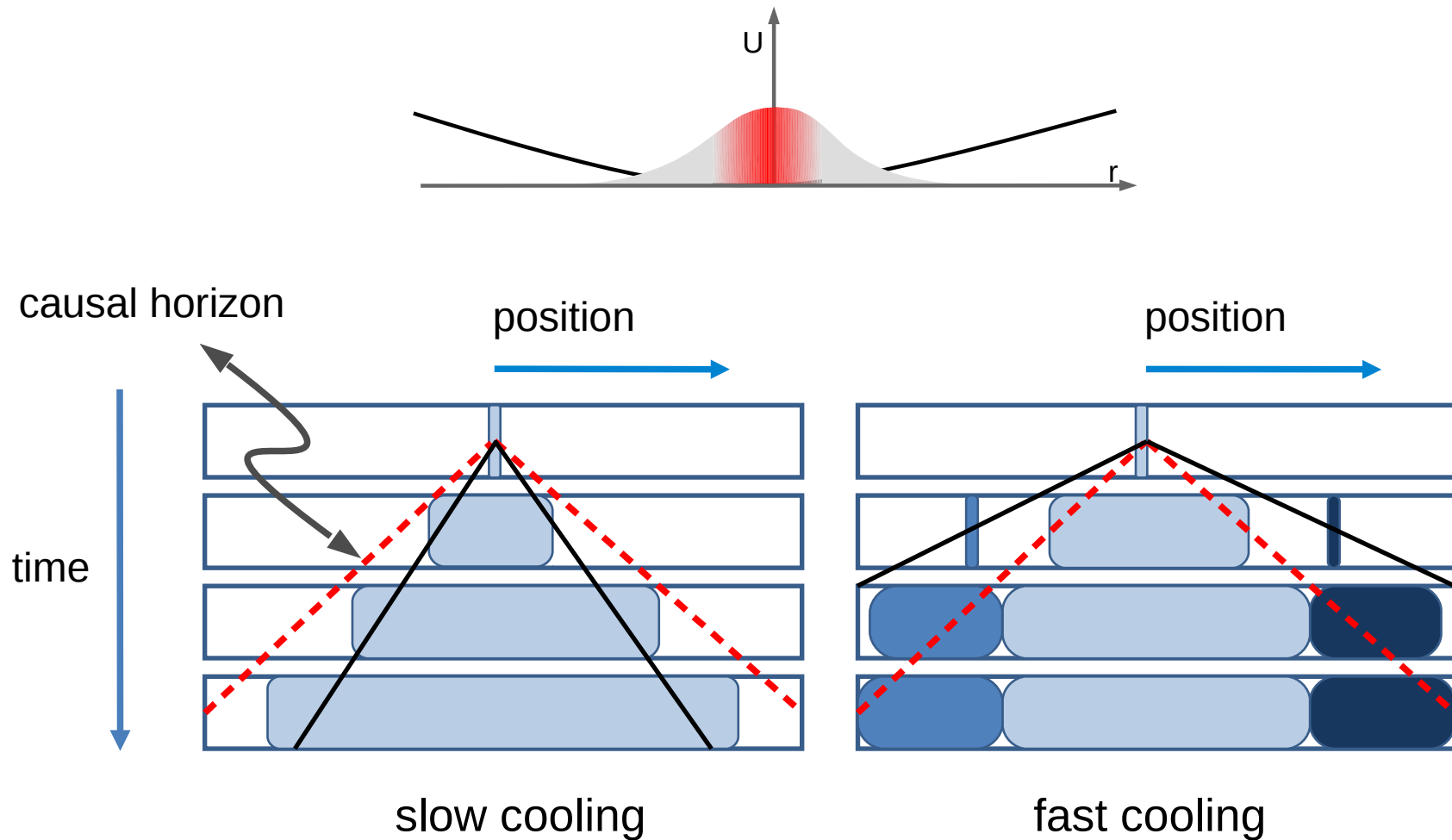
In trap

after expansion



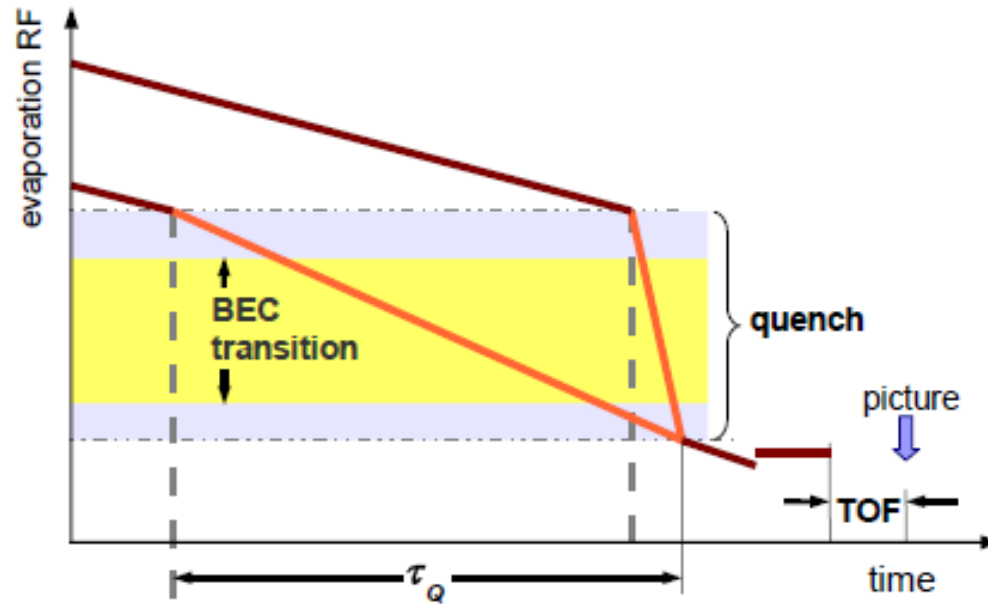
Key observation: the number of defects strongly depends on the rate at which the BEC transition is crossed !!

Generating solitons via the Kibble-Zurek mechanism

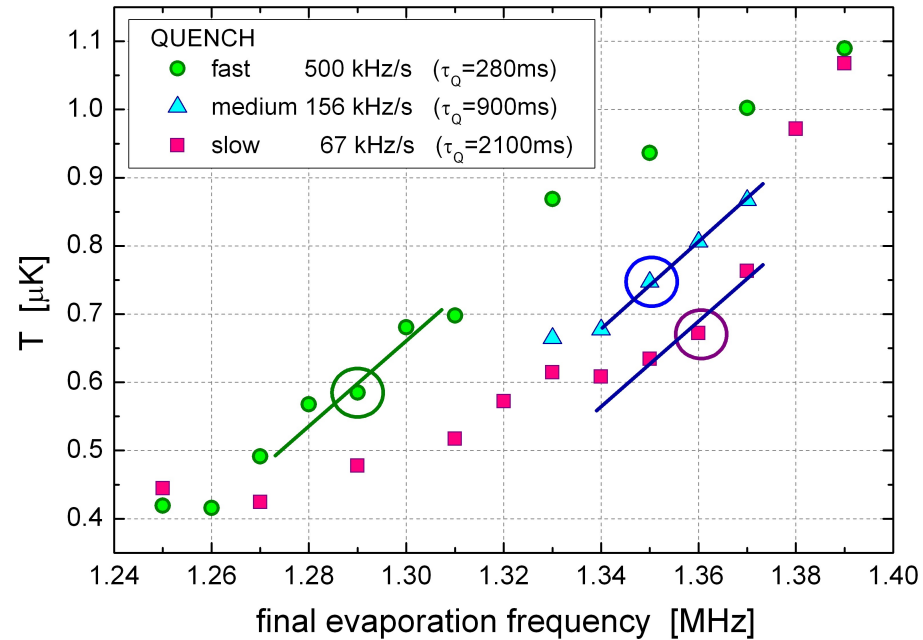


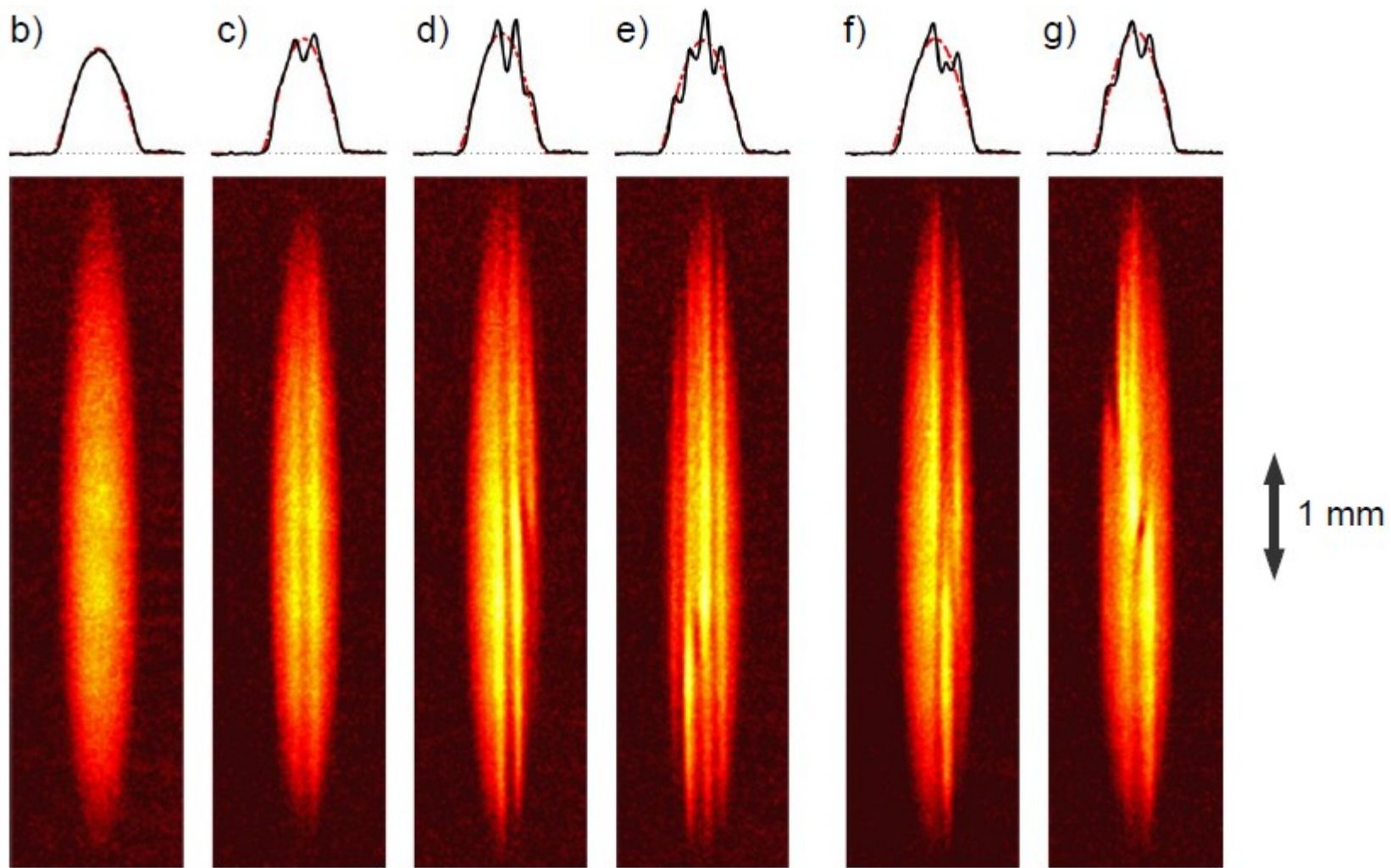
W. H. Zurek, PRL 102, 105702 (2009)

Check: change the cooling/quench time.

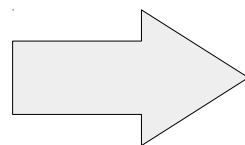


temperature VS evapoation thershold



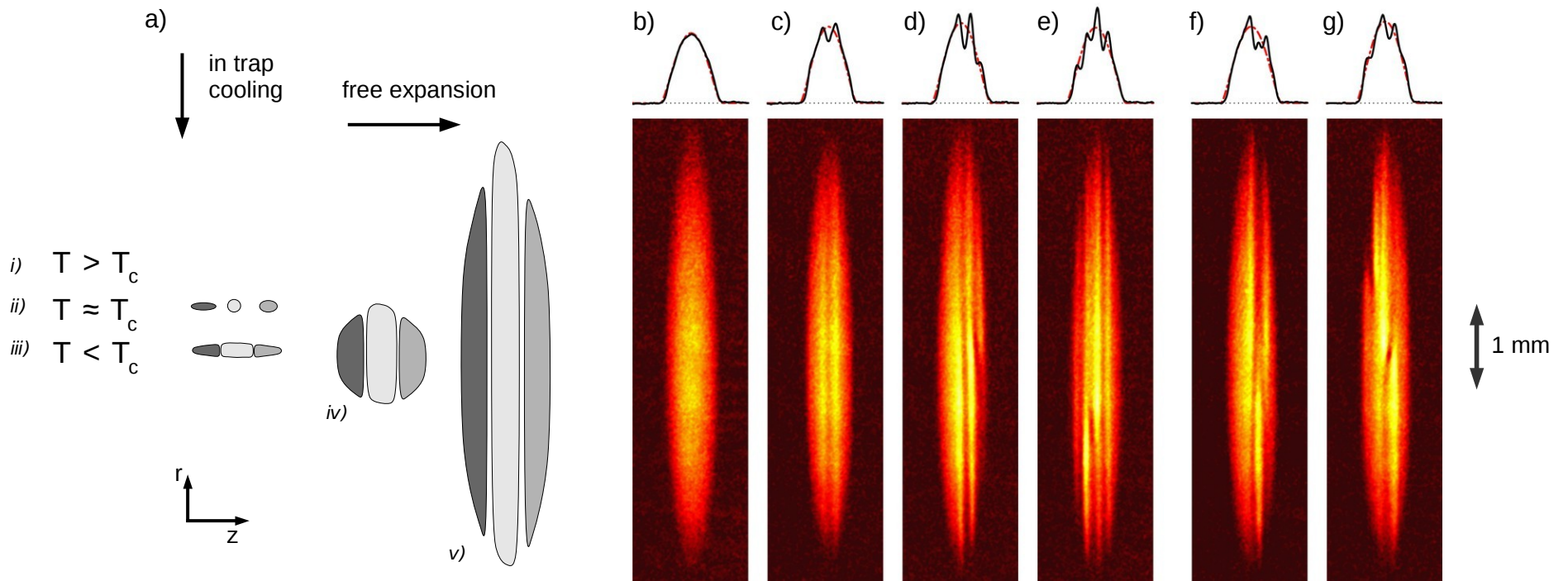


slow cooling

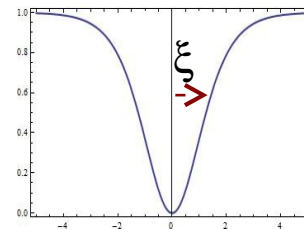


fast cooling

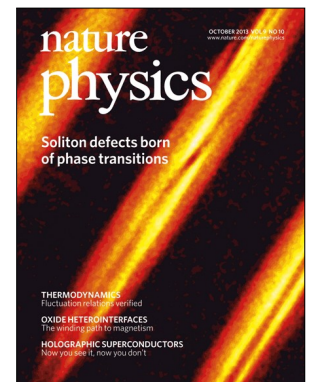
guess: these are gray solitons spontaneously nucleated at the BEC transition by the Kibble-Zurek mechanism (KZM) !!



imaging resolution: $10 \mu\text{m}$
 soliton width in trap: $\xi(0) = 200\text{-}250 \text{ nm}$
 width after TOF: $\xi(180 \text{ ms}) = 50\text{-}100 \mu\text{m}$



G.Lamporesi *et al.*, Nat. Phys. 9, 656 (2013)



Measurement of the KZ α coefficient

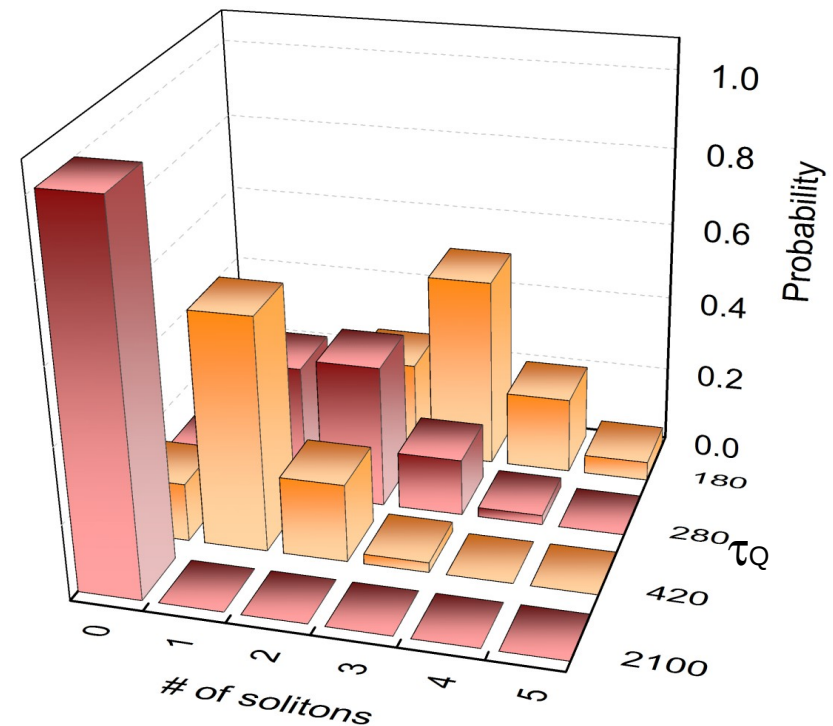
the number of defects is expected to follow a power-law as a function of the quench time (fixed size of the system)

$$N_s \propto \tau_Q^{-\alpha}$$

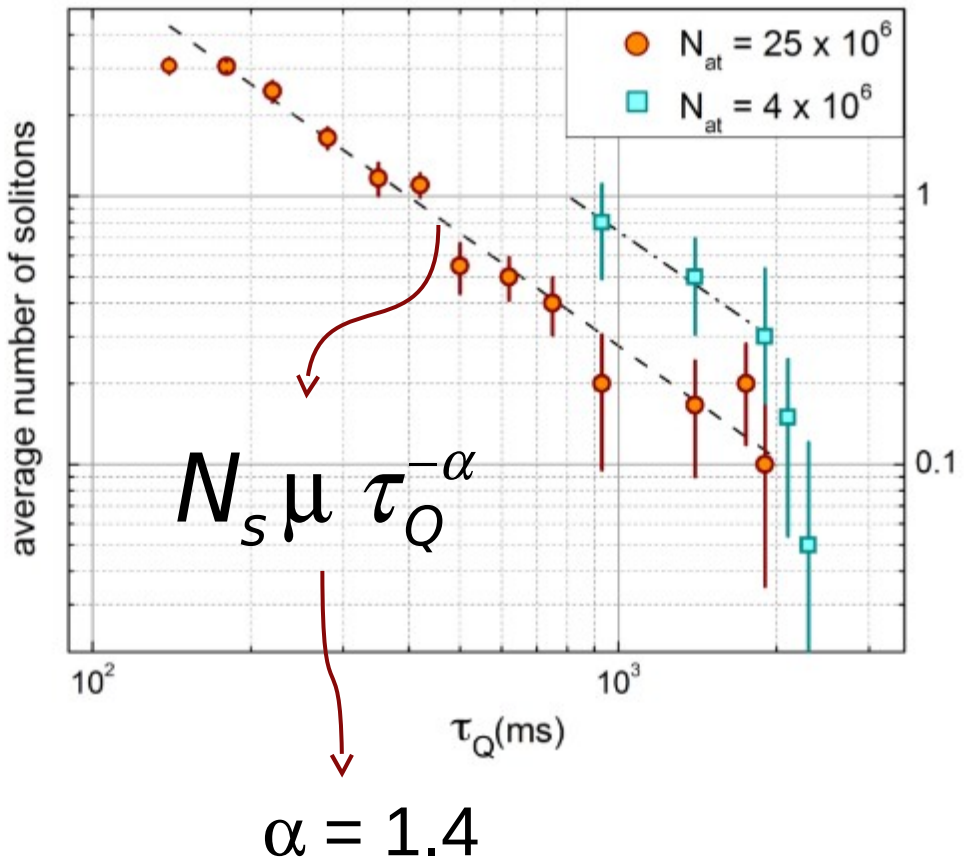
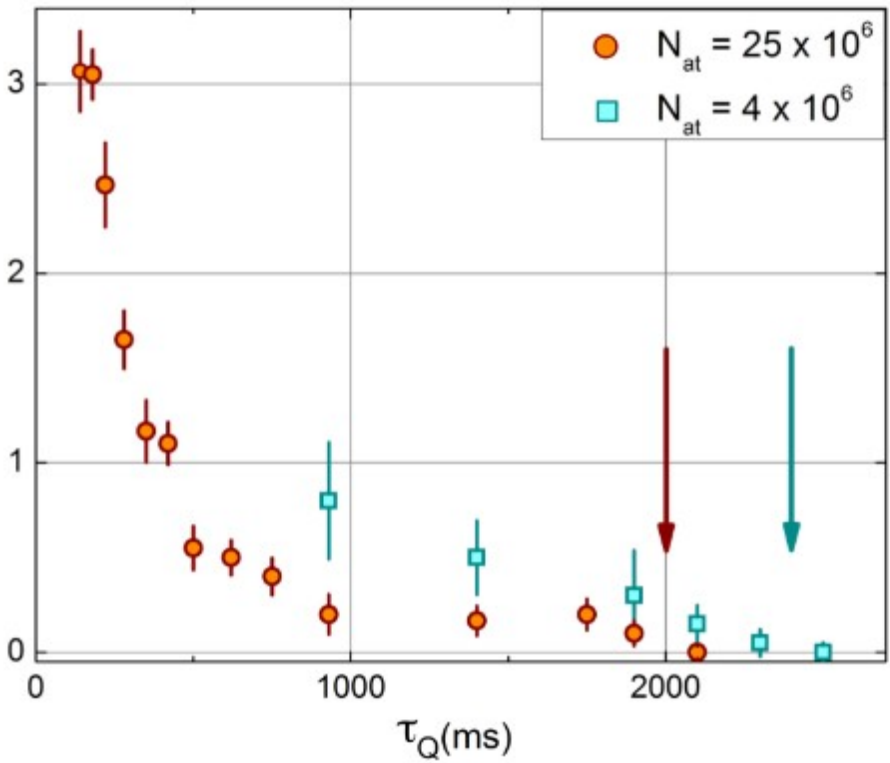
where α is determined by the critical exponents of the phase transition.

W. H. Zurek
PRL 102, 105702 (2009)

OK, we can count our solitons !



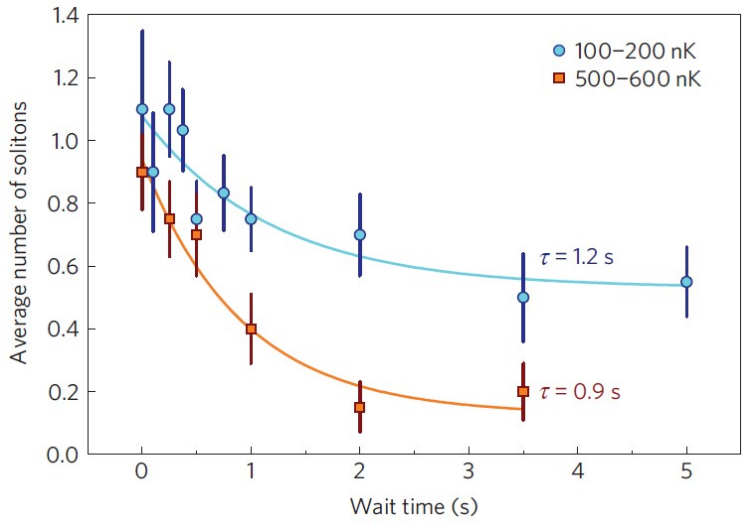
Measurement of the KZ α coefficient



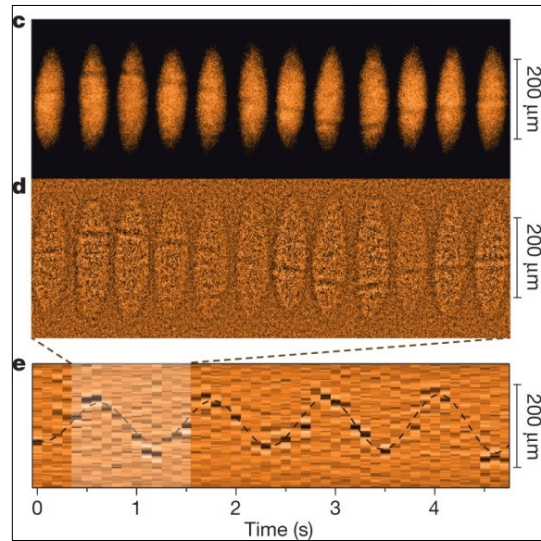
to compare with the available theoretical prediction (Zurek 2009)
 1D, homogeneous temperature

$$\alpha = 7/6 \sim 1.17$$

The lifetime puzzle



(also in DFG at MIT)

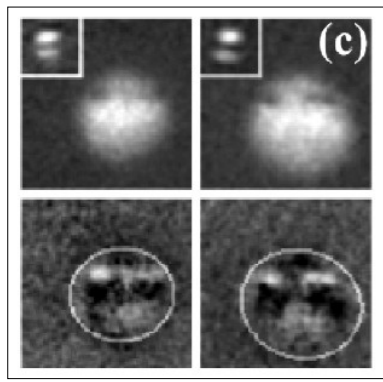
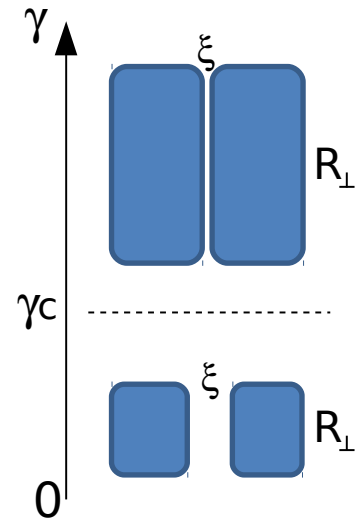


Yefsah *et al.*,
Nature **499**, 426 (2013)

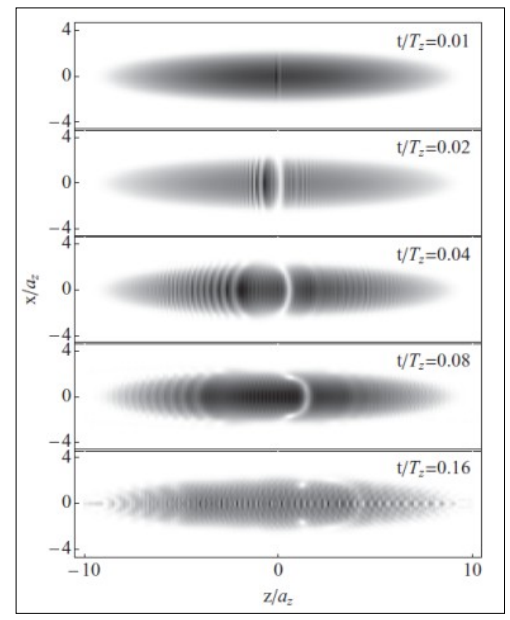
Solitons are expected to be **unstable**
THERMALLY (unless at T=0)
DYNAMICALLY (due to snake instabilities)

... and to decay into **vortex rings**

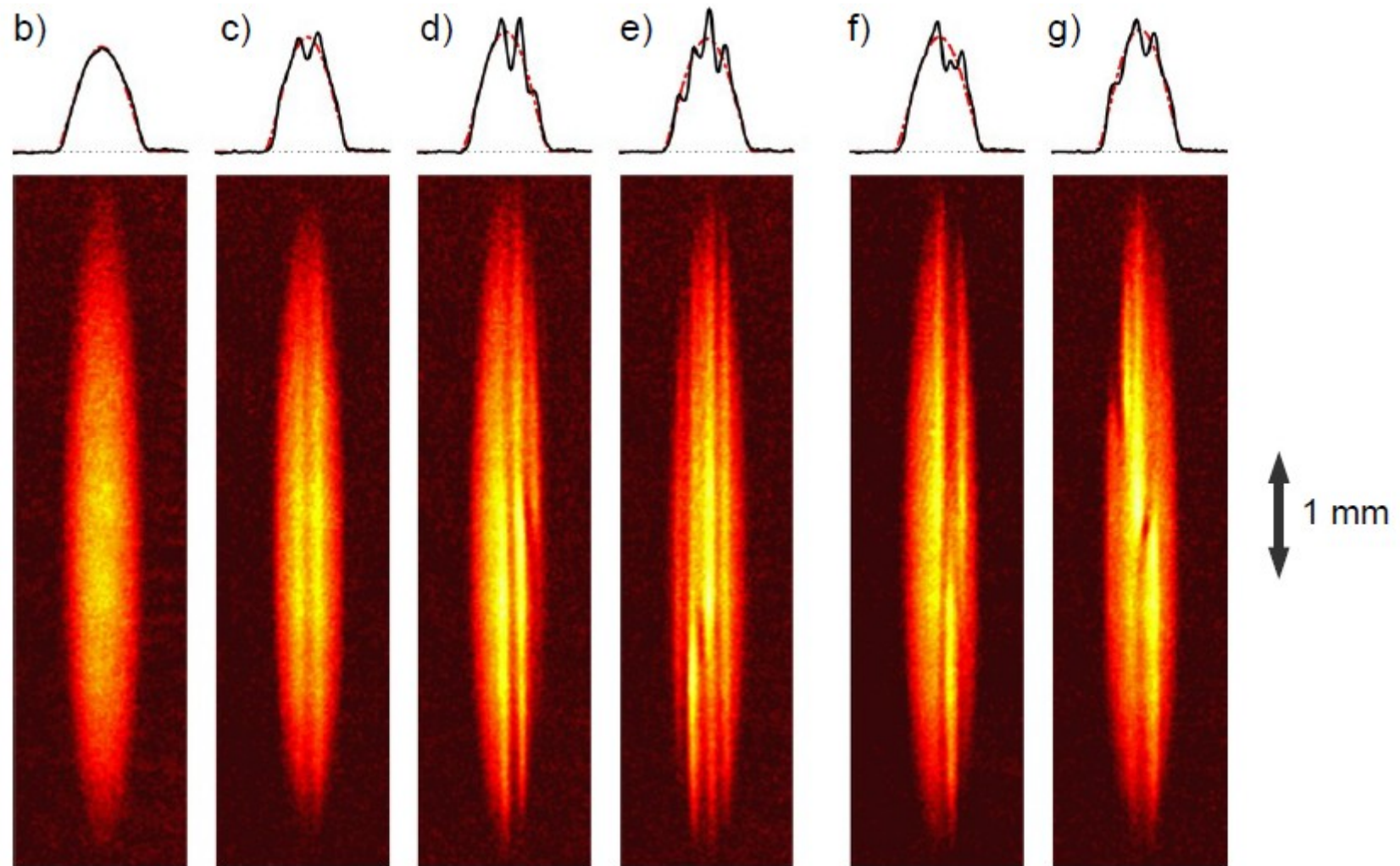
$$\gamma = \frac{\mu}{\hbar\omega_{\perp}} = \frac{R_{\perp}}{2\xi}$$



spherical BEC (JILA)
Anderson *et al.*,
PRL **86** 2926 (2001)



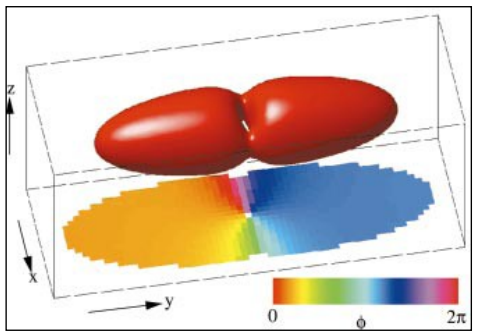
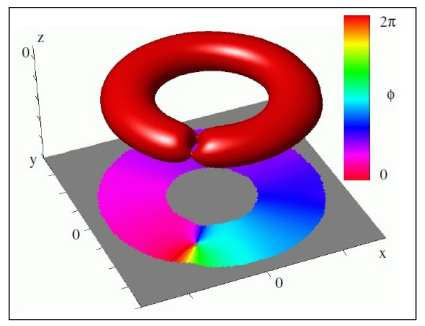
Reichl *et al.*,
PRA **88**, 053626 (2013)



Solitonic vortices

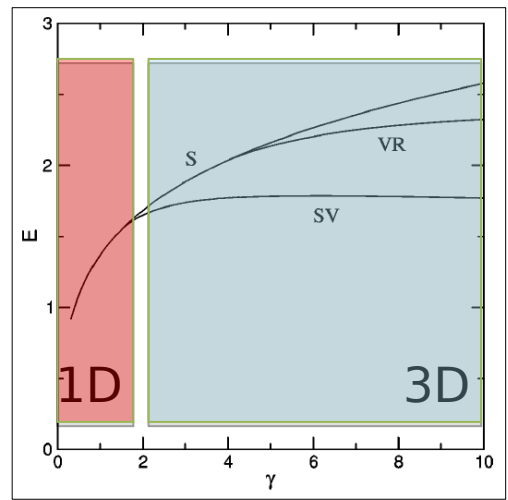
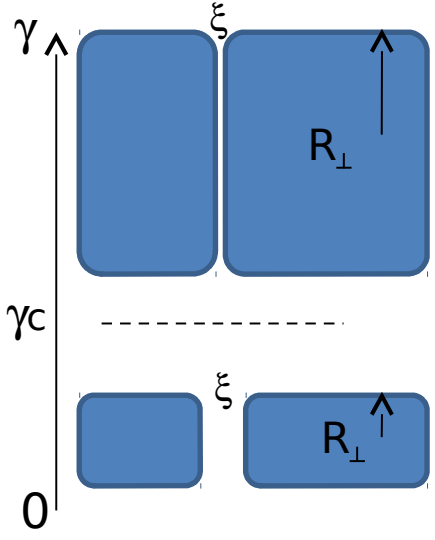
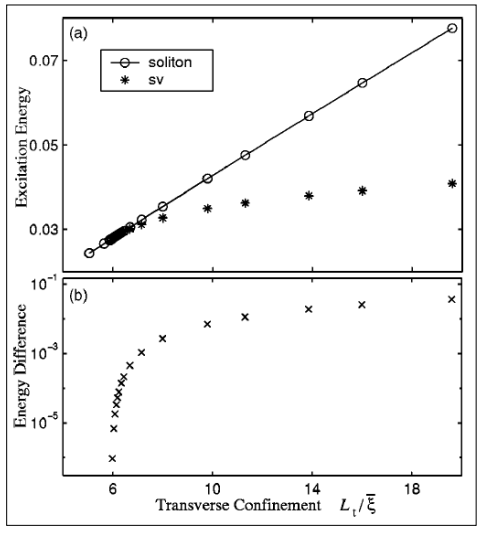
Vortex oriented perpendicularly to the axis of an axisymmetric elongated trap.

- Quantized vorticity
- Anisotropic phase pattern
- Planar density depletion



$$\gamma = \frac{\mu}{\hbar\omega_{\perp}} = \frac{R_{\perp}}{2\xi}$$

Brand *et al.*, JPB **34**, L113 (2001)



Komineas *et al.*, PRA **68**, 043617 (2003)

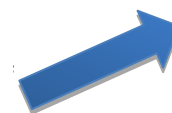
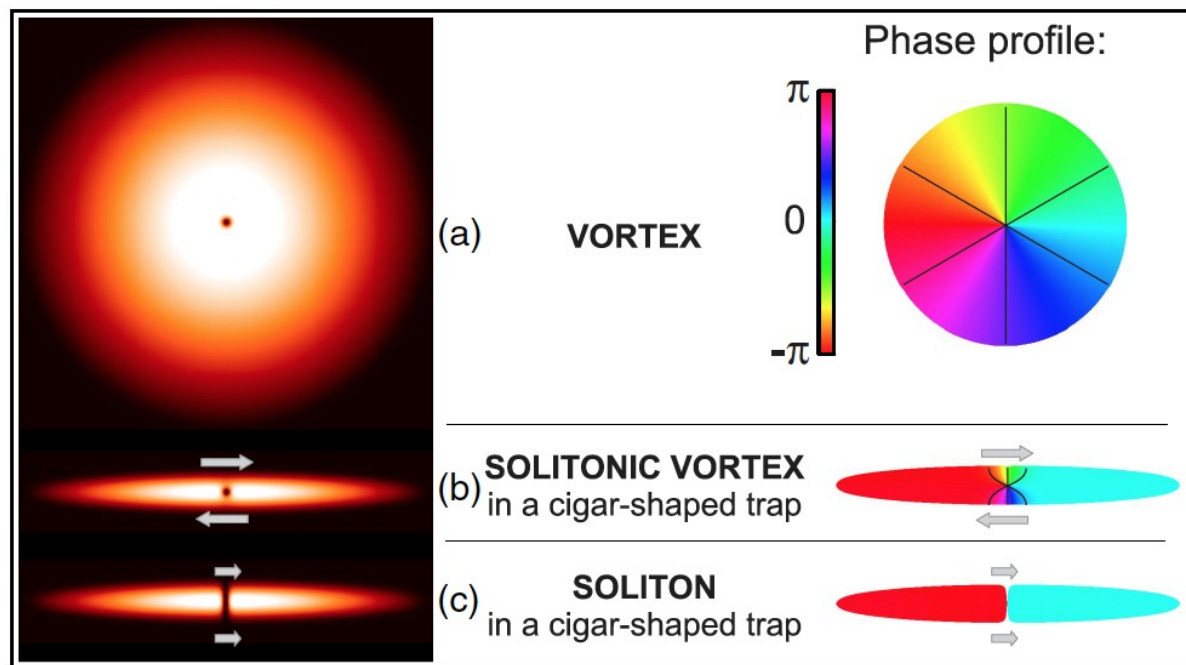
Brand *et al.*, PRA **65**, 043612 (2002)

Solitonic vortices

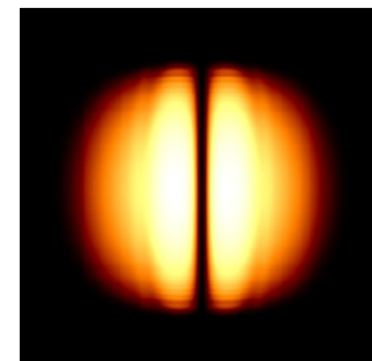
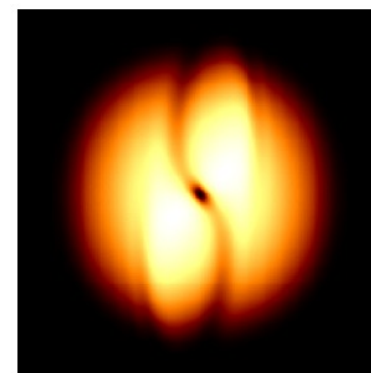
Density in trap

Phase

**Density
after free
expansion**

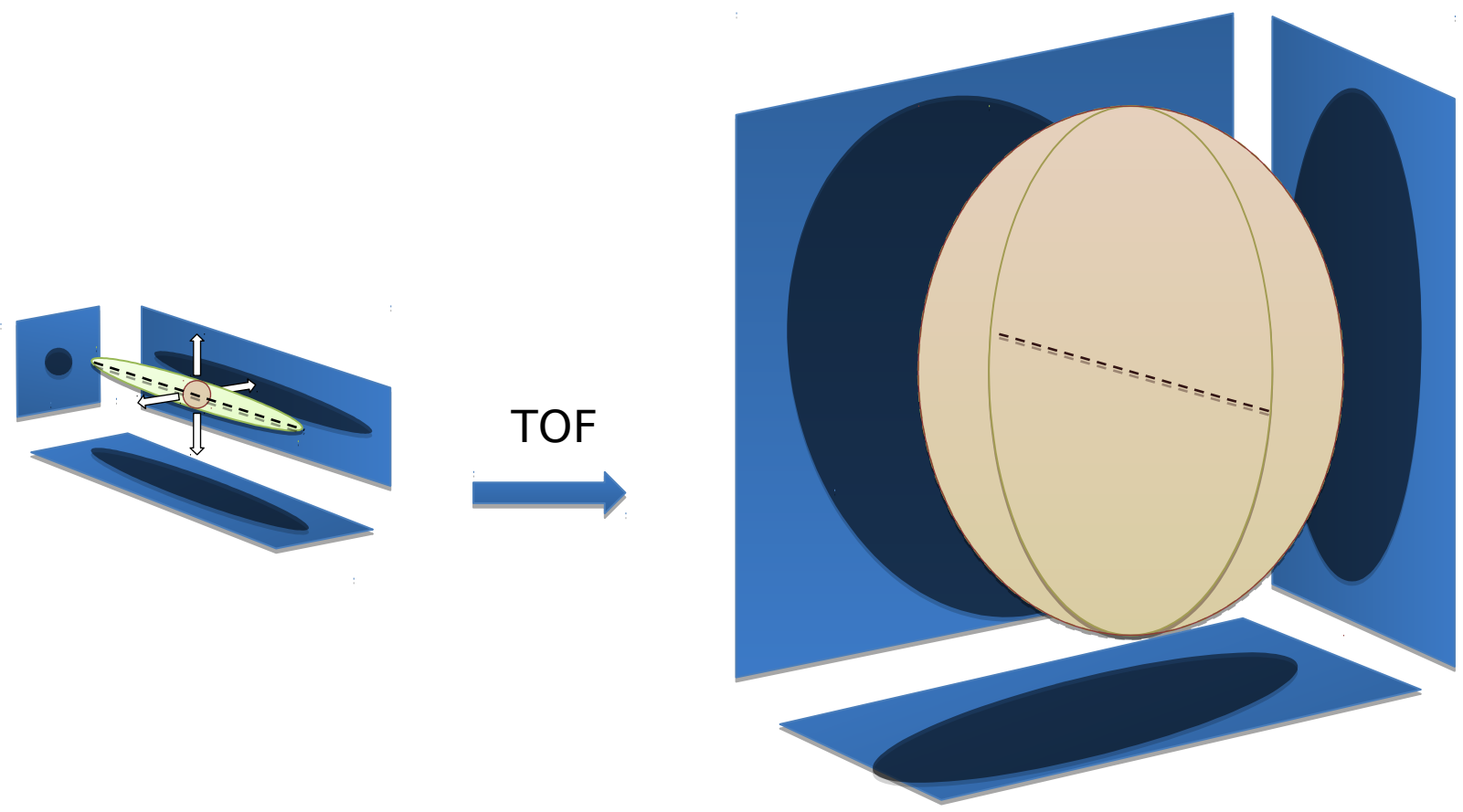


Asymmetric
twist



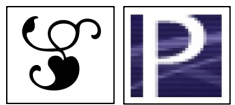
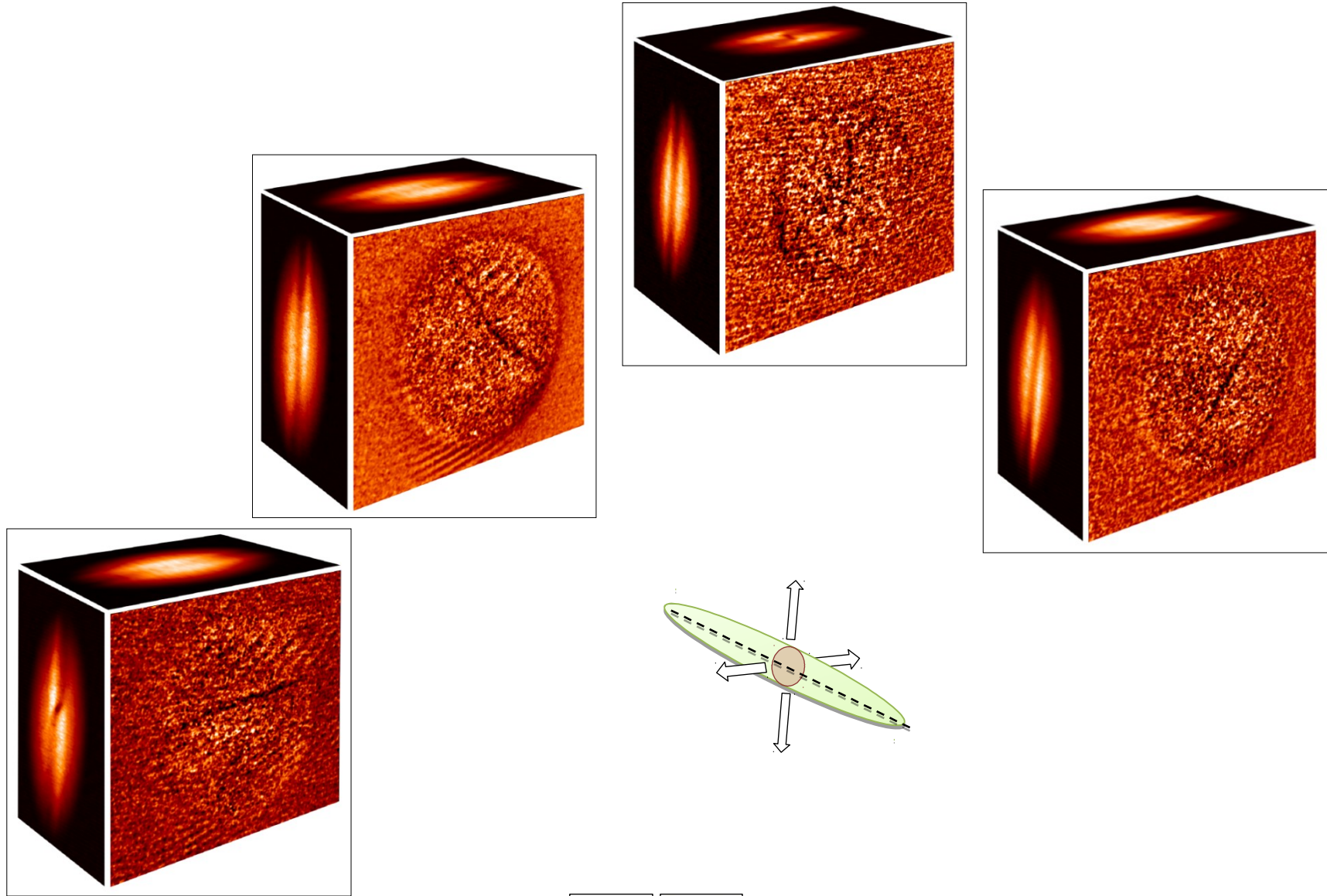
M. Tylutki *et al.*, EPJ-ST **224**, 577 (2015)

Solitonic vortices



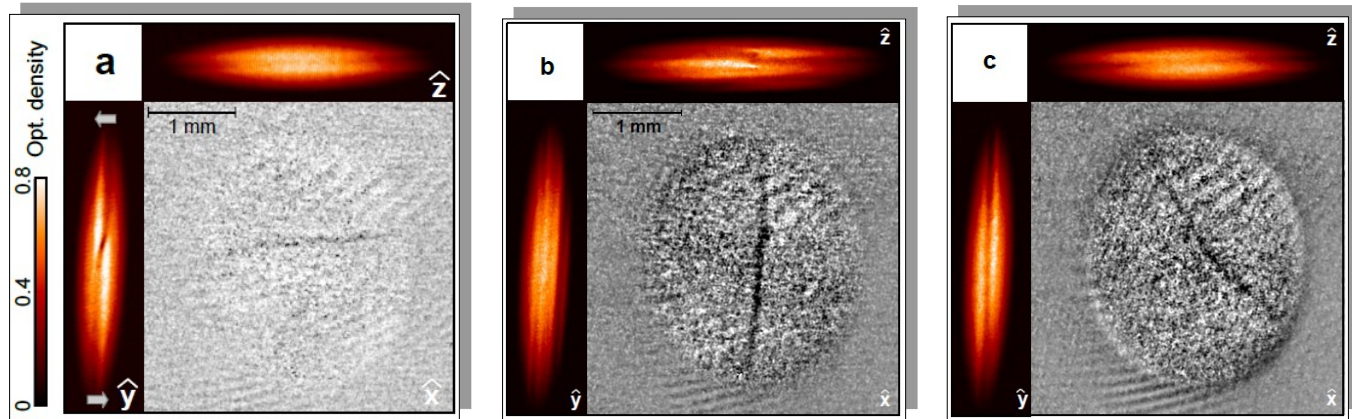
Random orientation

Triaxial absorption imaging after long TOF

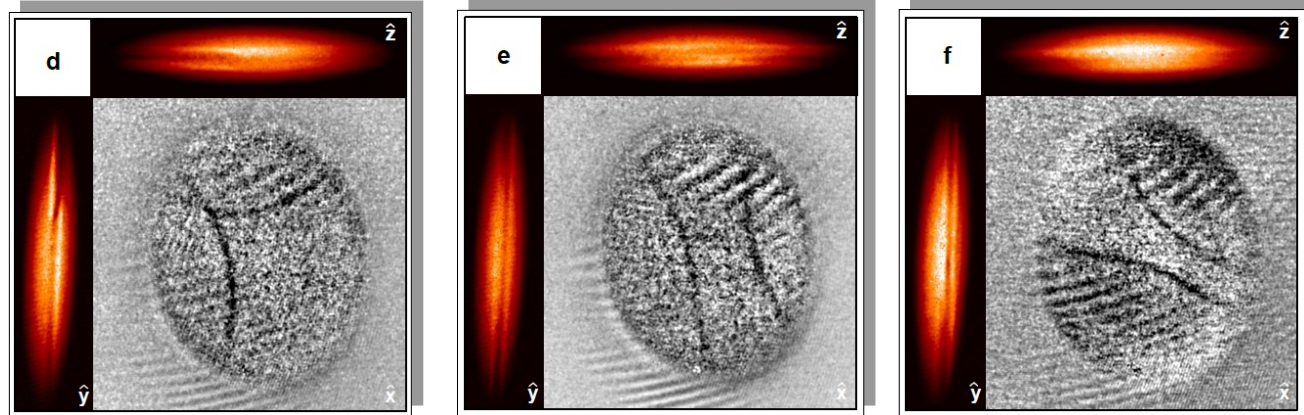


Random number

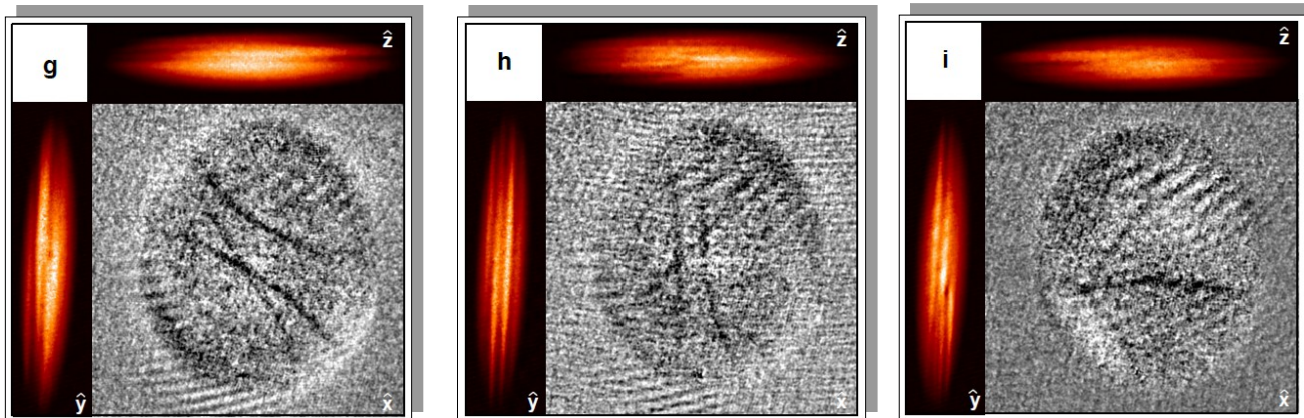
1



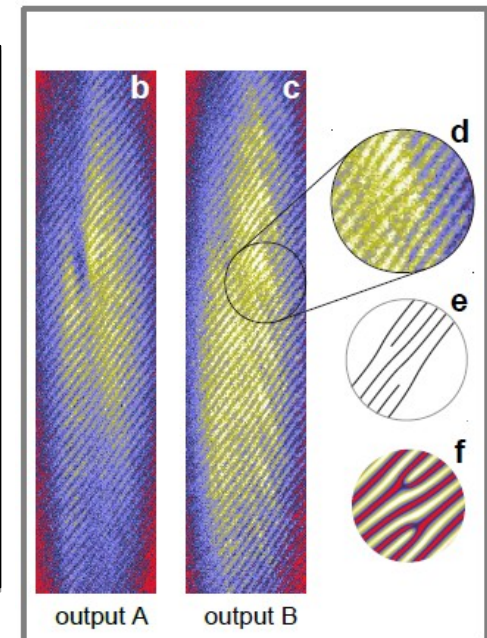
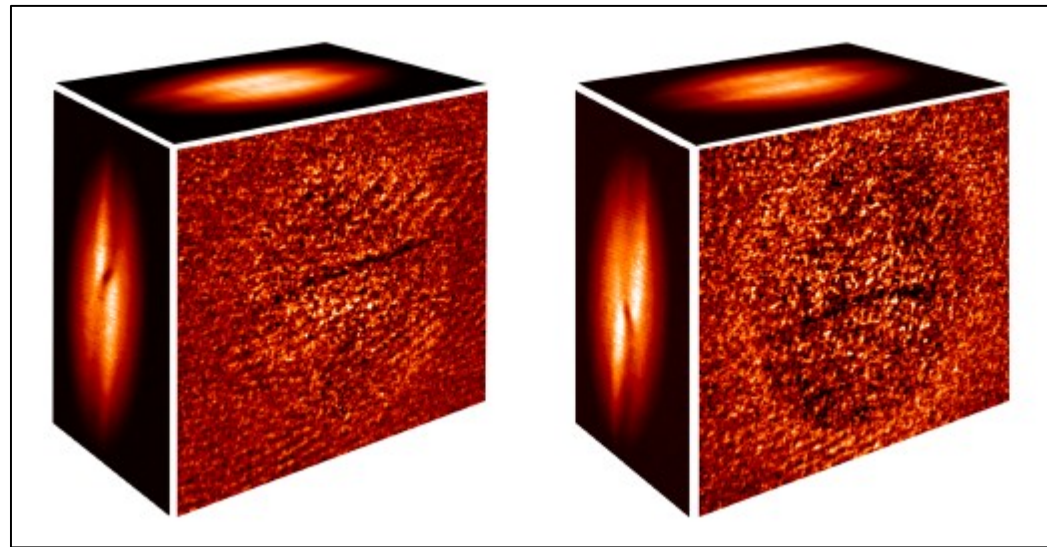
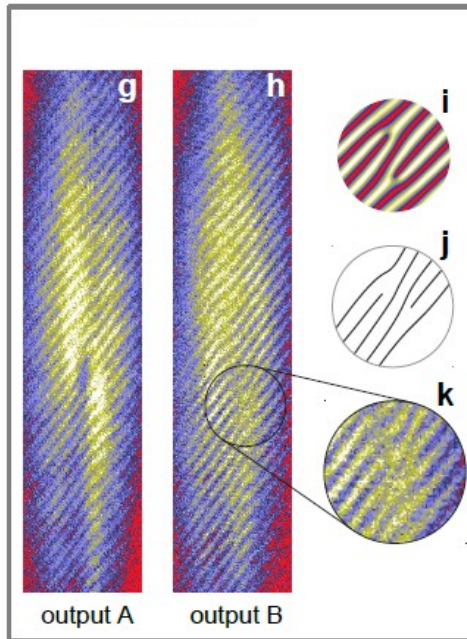
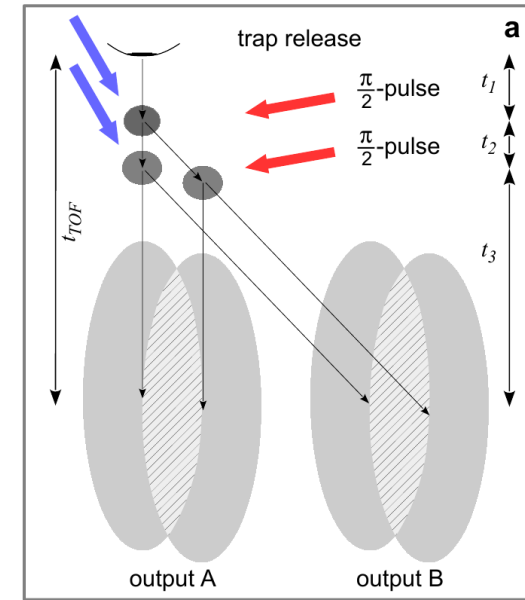
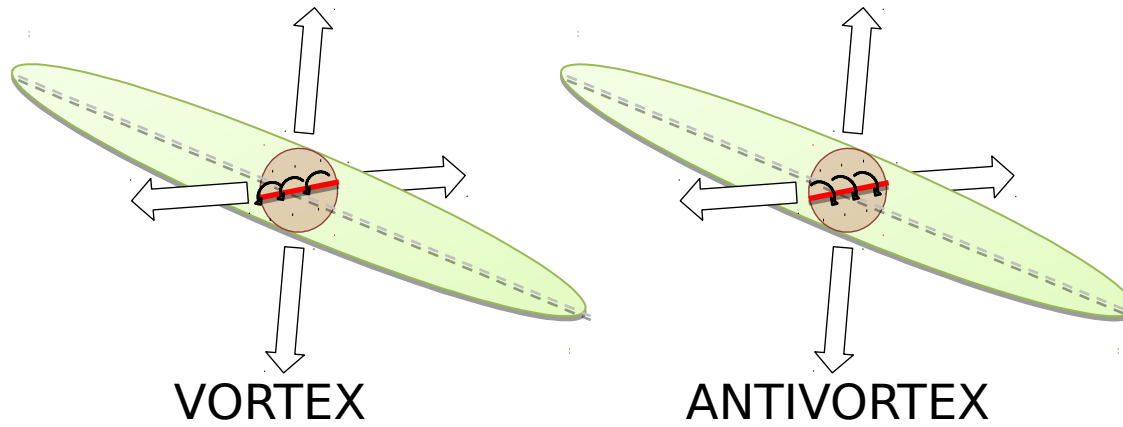
2



3



Random circulation

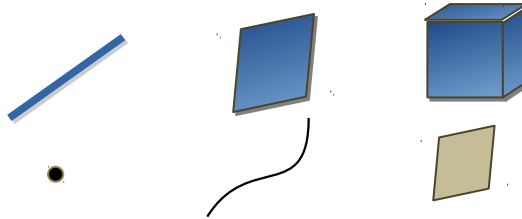


homodyne detection of the phase pattern by interfering two copies of the condensate

Scaling exponents

$$n \sim \frac{\hat{\xi}^d}{\hat{\xi}^D} = \frac{1}{\xi_0^{D-d}} \left(\frac{\tau_0}{\tau_Q} \right)^{(D-d) \frac{\nu}{1+z\nu}}$$

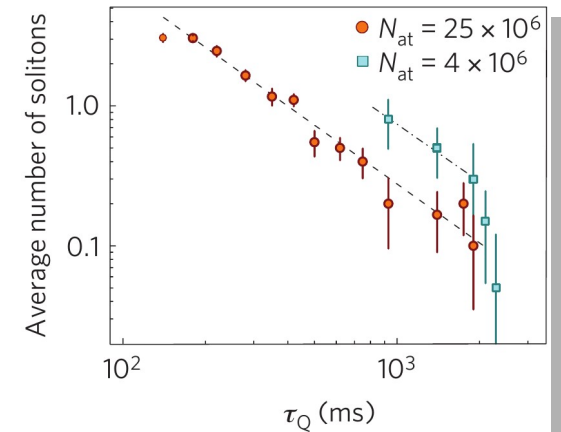
ν, z : critical exponents
 D : system dimension
 d : defect dimension



(D-d)=2

Critical exponents \ Trap	Homogeneous	Harmonic	Toroidal
Arbitrary (ν, z)	$\frac{2\nu}{1+\nu z}$	$\frac{2(1+2\nu)}{1+\nu z}$	$\frac{1+3\nu}{1+\nu z}$
Mean-field theory ($\nu = \frac{1}{2}, z = 2$)	$\frac{1}{2}$	2	$\frac{5}{4}$
Experiments/F model ($\nu = \frac{2}{3}, z = \frac{3}{2}$)	$\frac{2}{3}$	$\frac{7}{3}$	$\frac{3}{2}$

Del Campo *et al.*, NJP **13**, 083022 (2011)



($\alpha=1.4$)

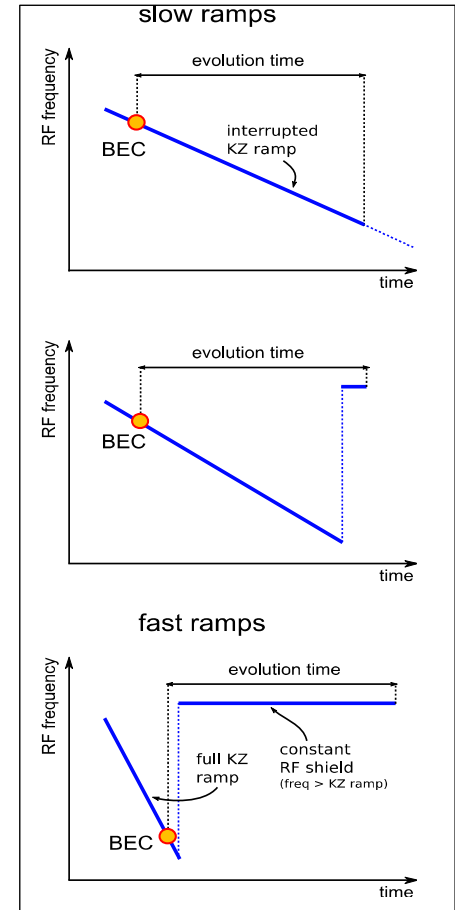
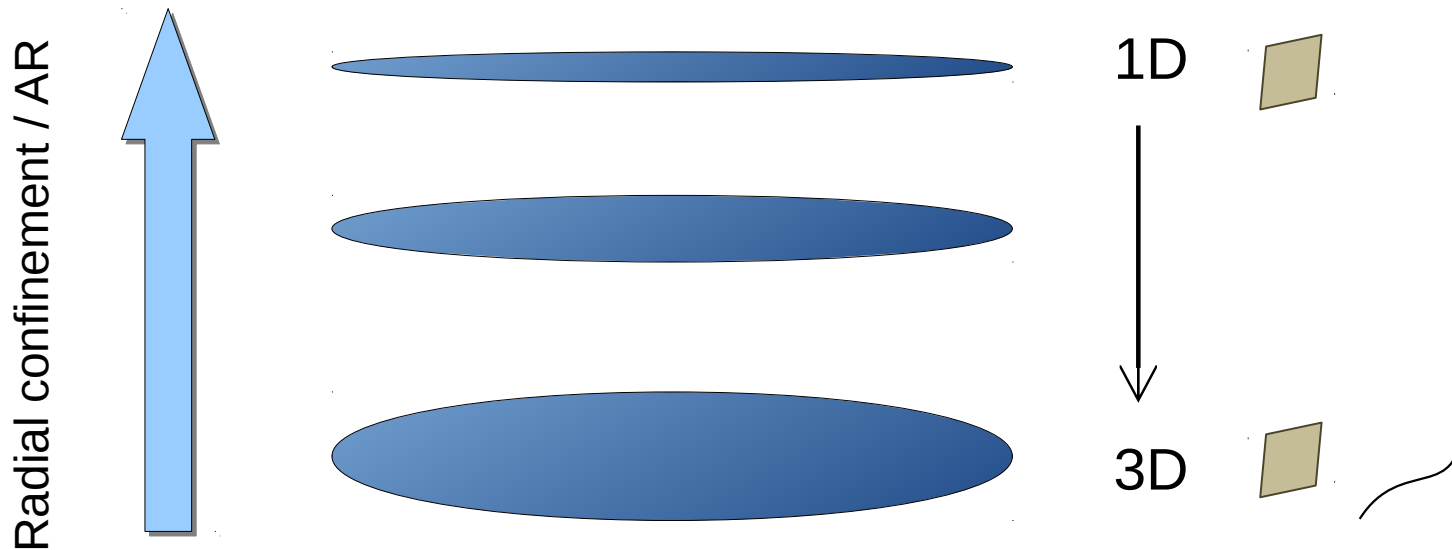
(D-d)=1

Homog.	Harm.
$\frac{\nu}{1+\nu z}$	$\frac{1+2\nu}{1+\nu z}$
1/4	1
1/3	7/6

Zurek, PRL **102**, 105702 (2009)

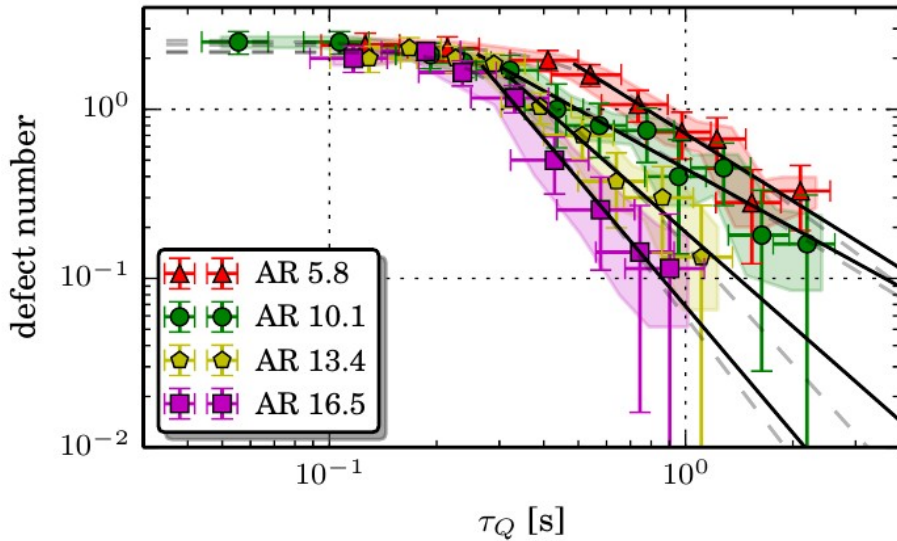
Search for effects on KZ scaling due to geometry of confinement:

- dimensional cross-over
- creation of different types of defects

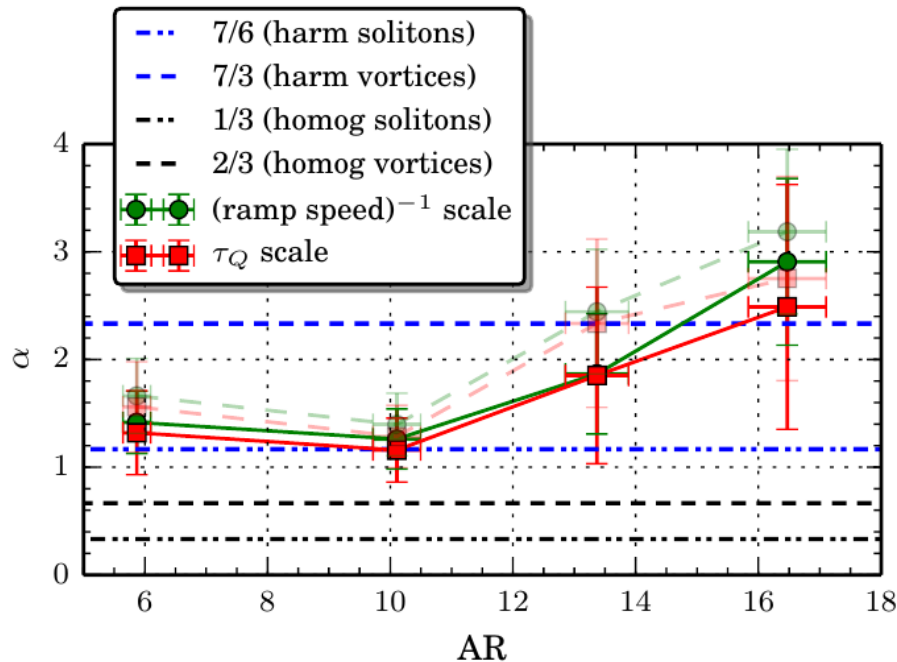


Revised evaporation ramps suppressing the effects of decay of defects

preliminary



- power-law scaling for slow ramps
- aspect ratio dependent exponent



- flat plateau for fast ramps
- plateau independent on aspect ratio

	model	α_{homog}	α_{harm}
solitons in 3D or in 1D ($S - D = 1$)	MF	1/4	1
	F-model	1/3	7/6
vortices in 3D ($S - D = 2$)	MF	1/2	2
	F-model	2/3	7/3

Dynamics of quantized vortices

Determine dissipative and transport processes in:

- Superfluid helium
- Superconductors
- Neutron stars

In atomic BECs:

Controllable environment, spatial scale from ξ to tens of ξ , inhomogeneous systems, boundary physics...

BUT

Vortices are produced stochastically and their dynamics cannot be followed through standard destructive absorption imaging

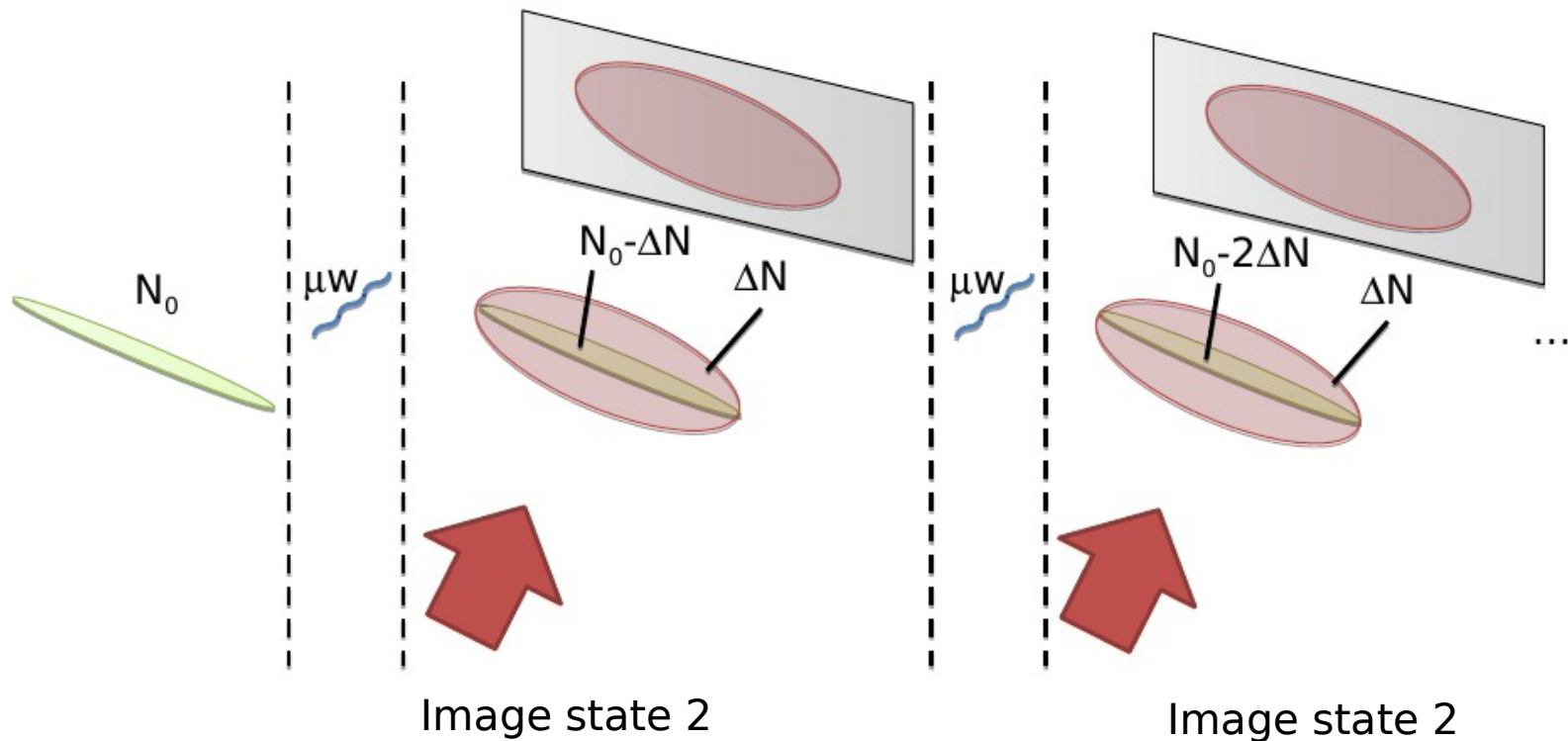
Stroboscopic imaging of defect dynamics

μ -wave pulses extract a small fraction from the BEC

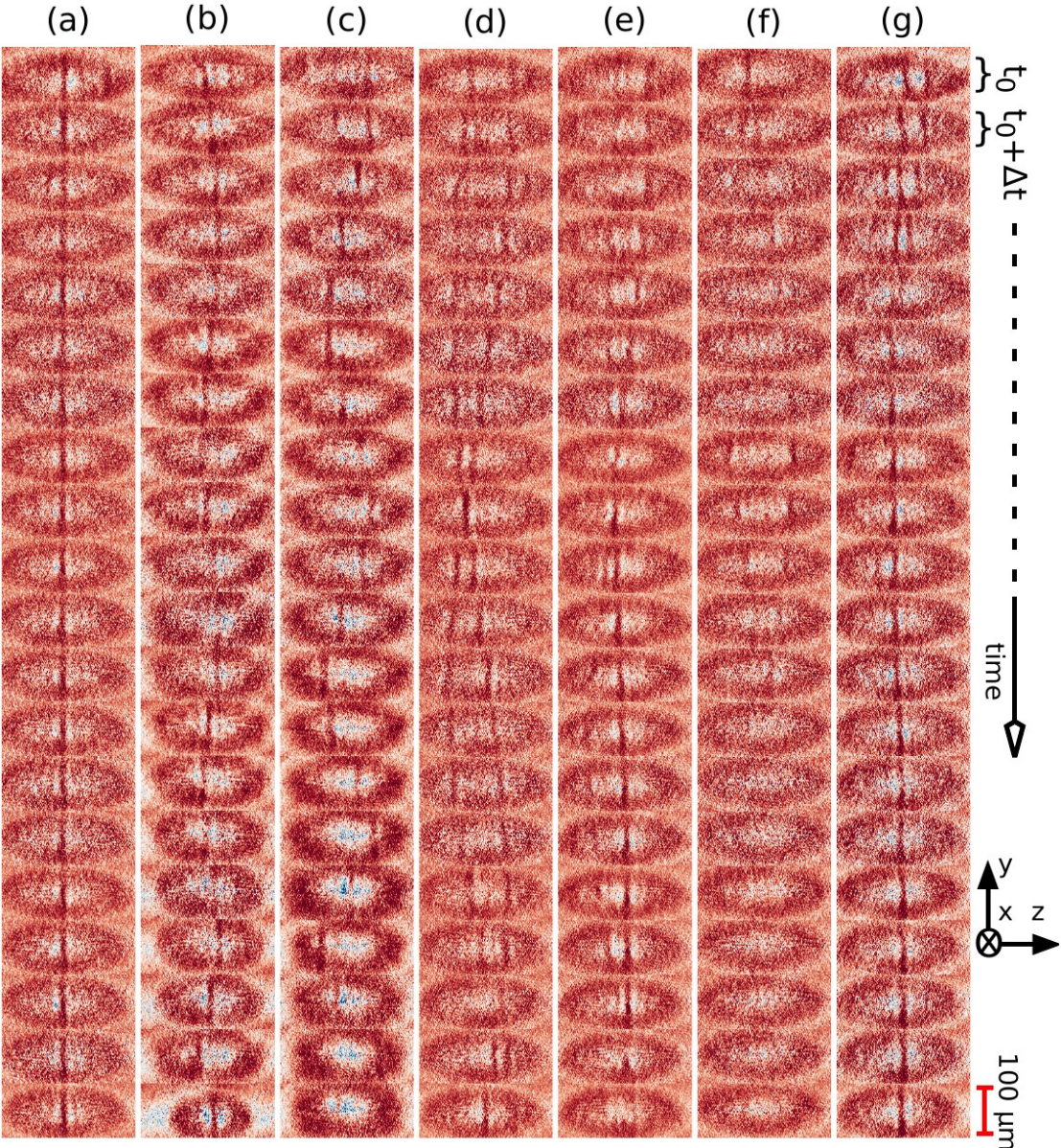
$$|1, -1\rangle \longrightarrow |2, -2\rangle$$

$$\frac{\Delta N}{N_0} \sim 4\%$$

- Initial atom number $\sim 10^7$
- Magnetic harmonic trap in $|1, -1\rangle$ with $\{\omega_{x,y} = \omega_{\perp}, \omega_z\}/2\pi = \{131, 13\}$ Hz
- 13 ms expansion in $|2, -2\rangle$ plus RF dressing
- Selective imaging of the output coupled fraction

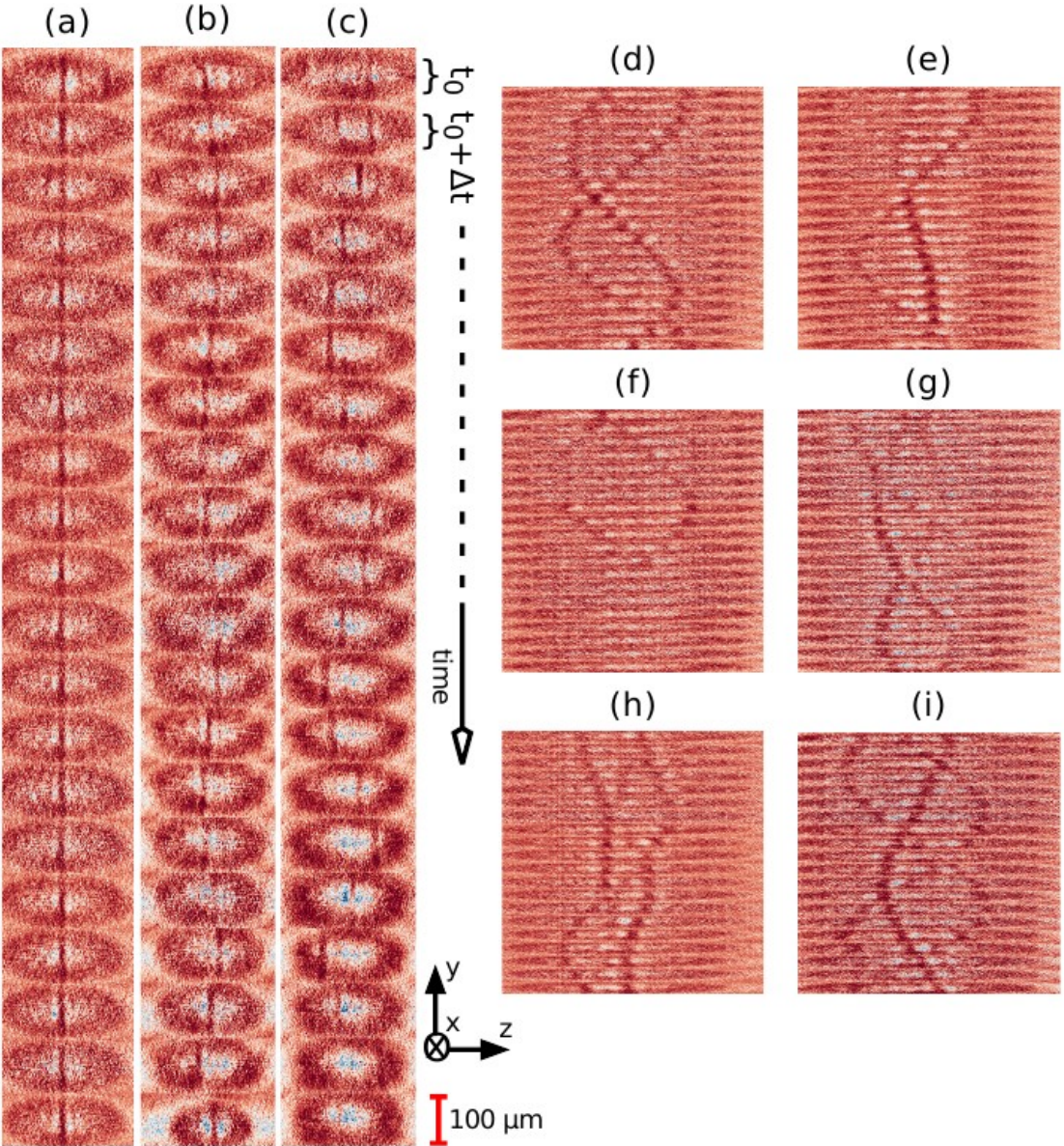


Stroboscopic imaging of defect dynamics



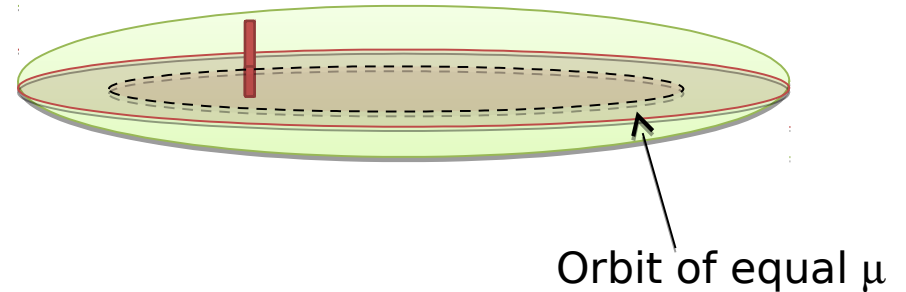
- expansion in the anti-trapped state
- selective imaging of the output coupled fraction
- imaging iterated up to 20 times
- plot residuals from the Thomas-Fermi profile

Stroboscopic imaging of defect dynamics



- expansion in the anti-trapped state
- selective imaging of the output coupled fraction
- imaging iterated up to 20 times
- plot residuals from the Thomas-Fermi profile

Vortex dynamics



SINGLE VORTEX DYNAMICS

A straight vortex line is expected to precess in an inhomogeneous non-rotating condensate, following an equipotential elliptical orbit around the center:

$$T_{SV} = \frac{4(1 - r_0^2)\mu}{3\hbar\omega_{\perp} \ln(R_{\perp}/\xi)} T_z$$

$$T_z = \frac{2\pi}{\omega_z} \quad \text{axial trapping period}$$

$$r_0 = \frac{z_{\max}}{R_z} = \frac{y_{\max}}{R_{\perp}} \quad \text{normalized oscillation amplitude}$$

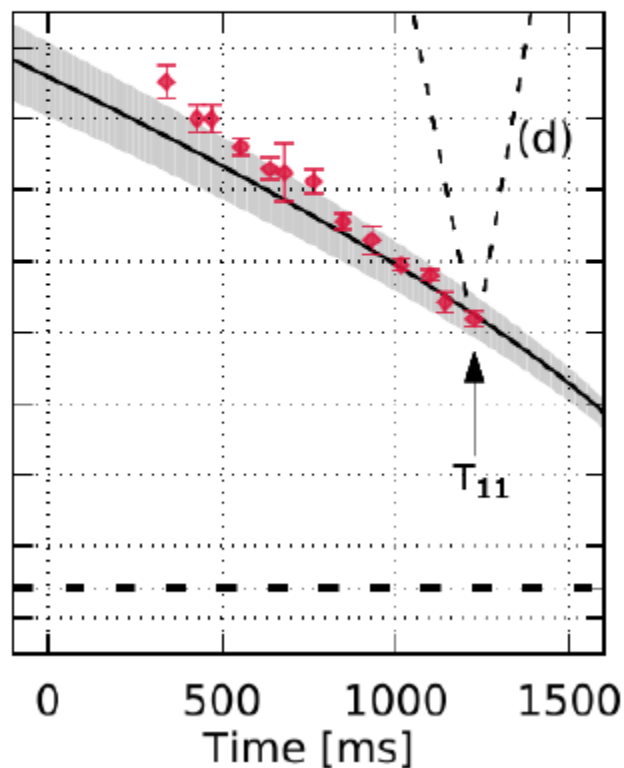
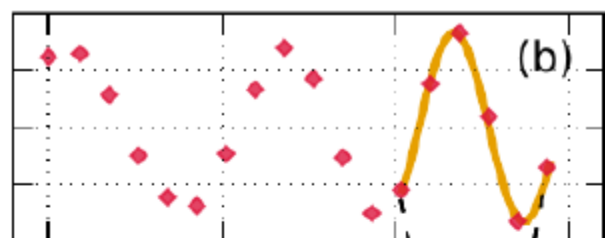
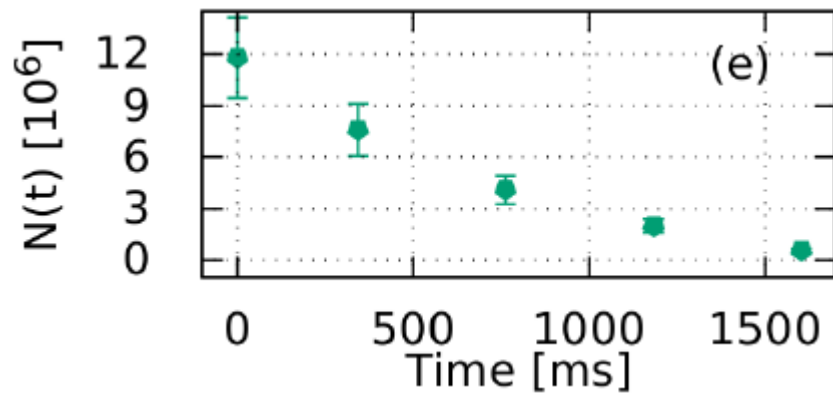
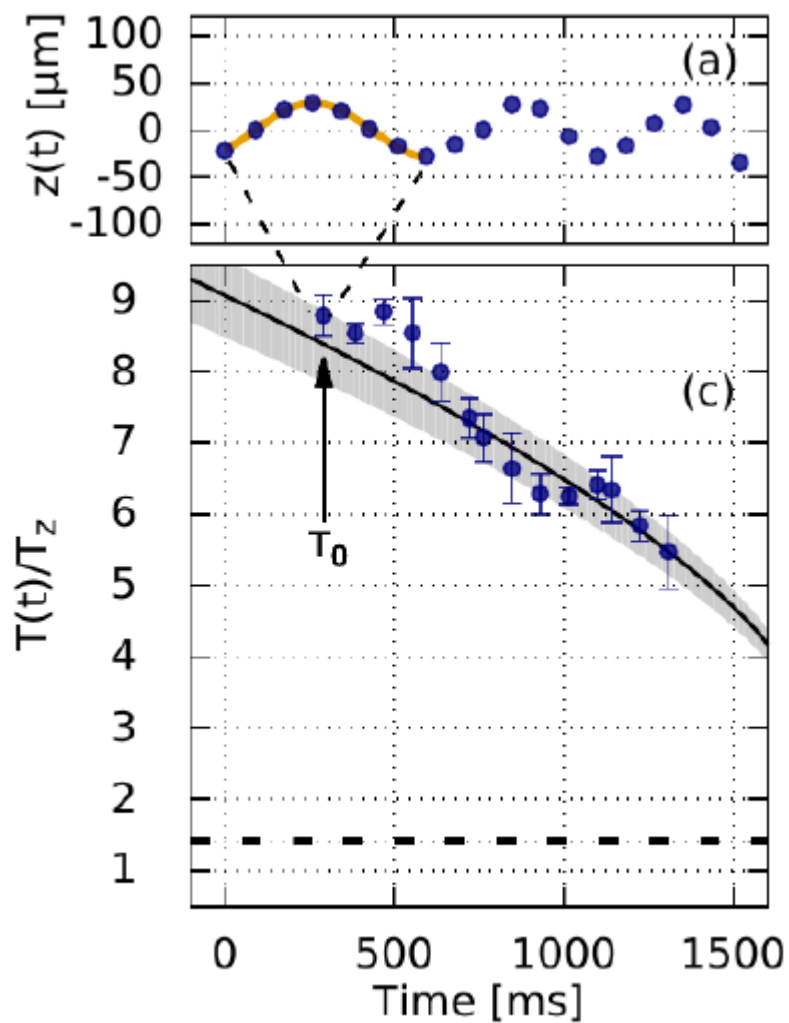
$$\xi \quad \text{condensate healing length}$$

A. L. Fetter and J.-K. Kim, J. Low Temp. Phys. **125**, 239 (2001)

L. P. Pitaevskii, arXiv: 1311.4693 (2013), M. J. H. Ku *et al.*, PRL **113**, 065301 (2014)

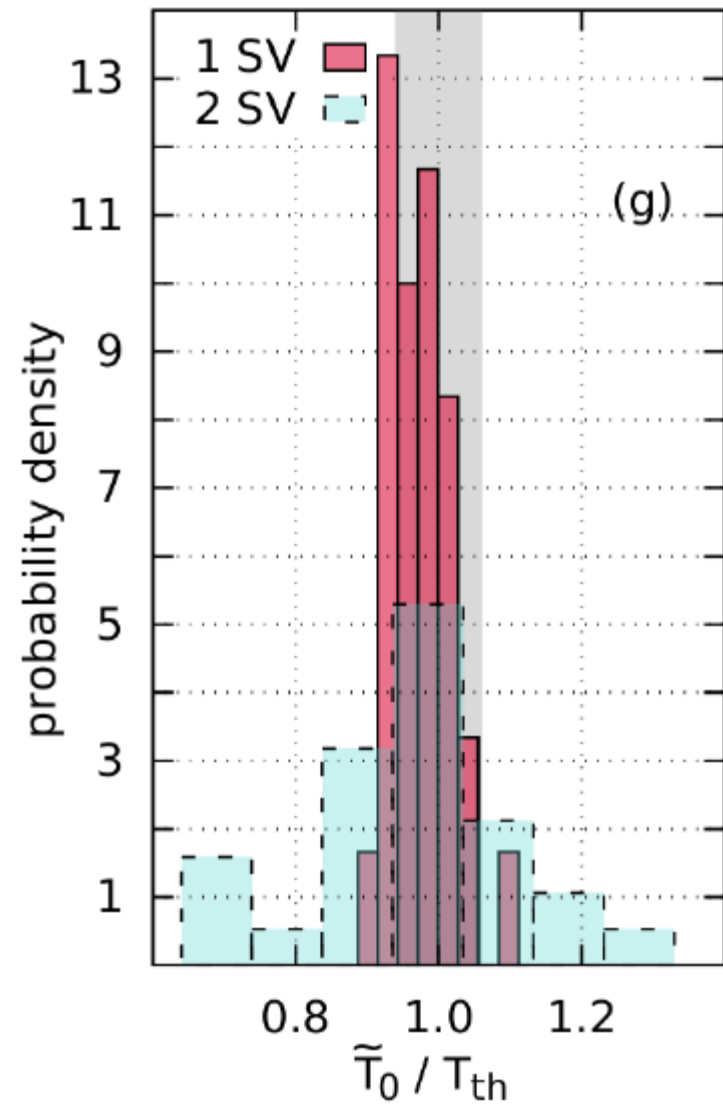
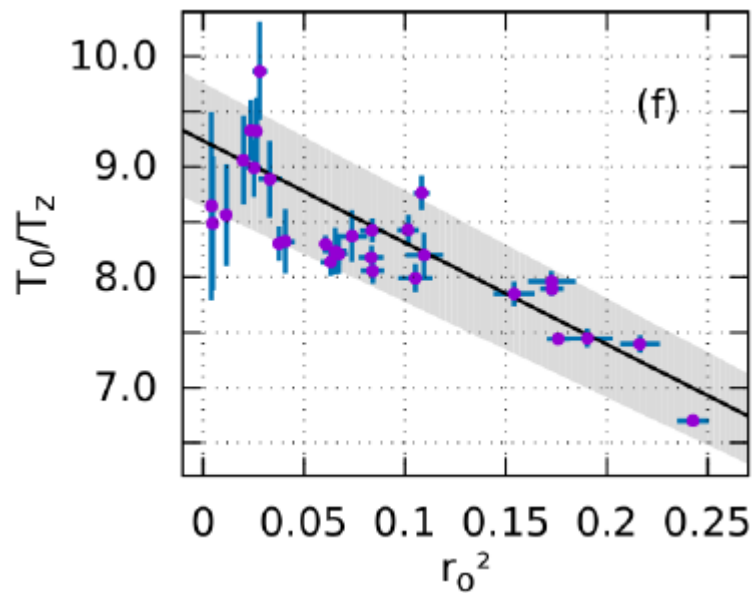
Period VS atom number

$$T_{SV} \propto \mu \propto N(t)^{2/5}$$



Period VS amplitude of orbit

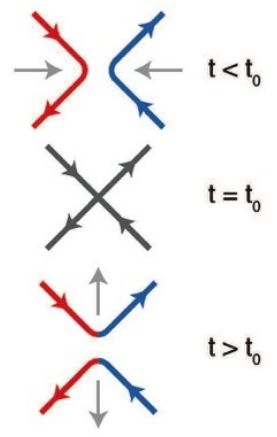
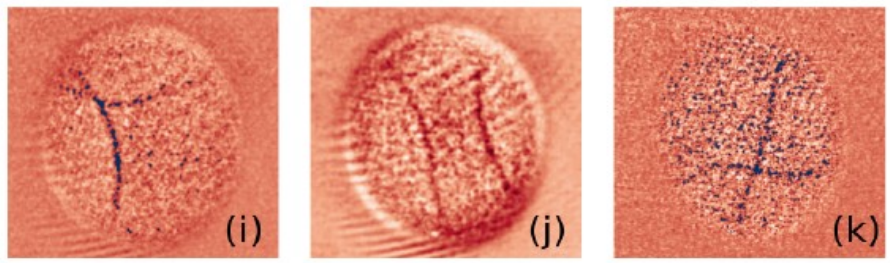
$$\frac{T_{SV}}{T_z} = \frac{4(1 - r_o^2)\mu}{3\hbar\omega_{\perp} \ln(R_{\perp}/\xi)}$$



Interaction among vortices

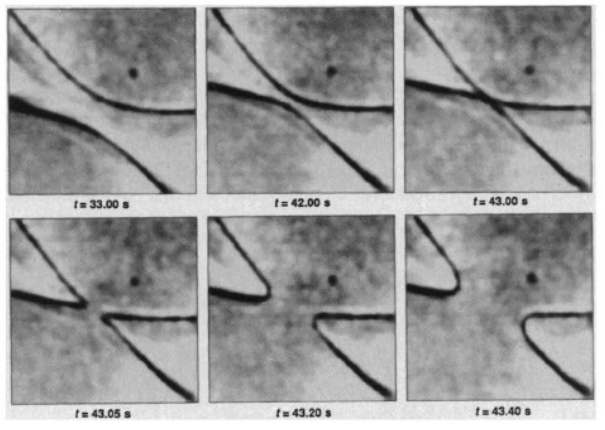
Random orientation of the nodal lines in the radial plane

→ Full 3D vortex interaction



- Ideal benchmark for:
- Vortex annihilation
 - Vortex decay
 - Vortex reconnection

Reconnection in liquid crystals



Present simulations:
Vortices are initially at rest

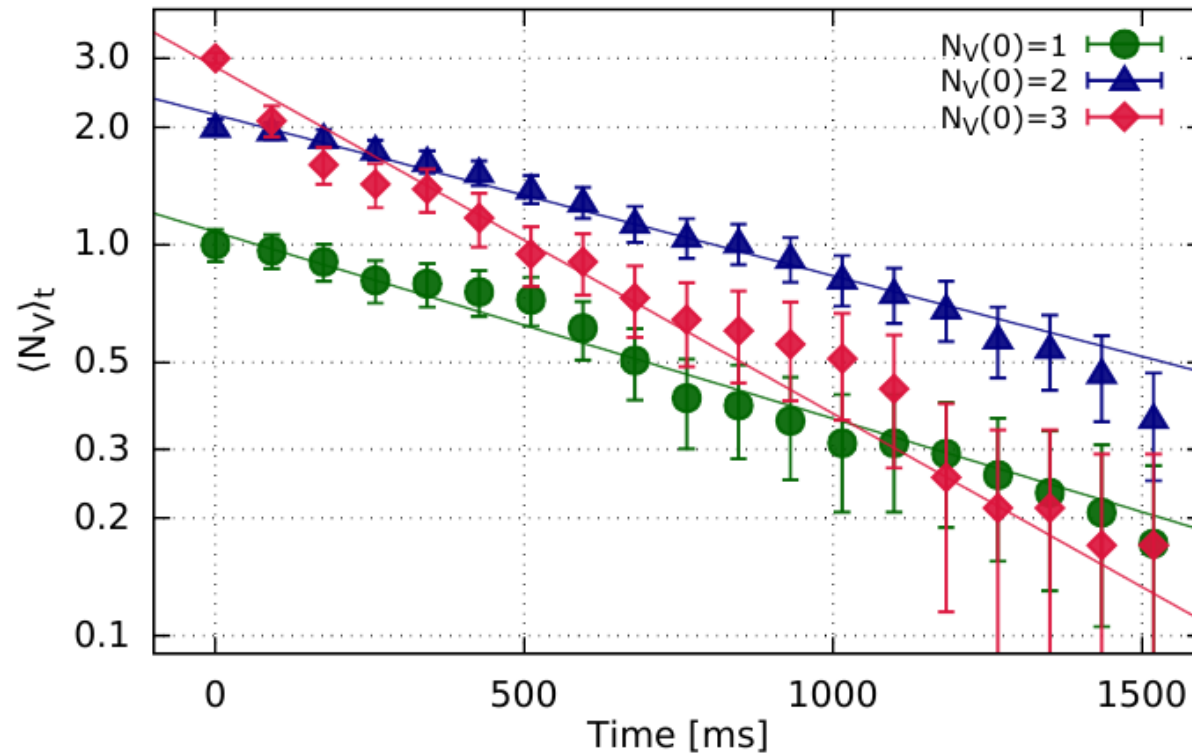
Our experiment:
finite relative momentum

Chuang *et al.*, Science **251**, 1336 (1991)

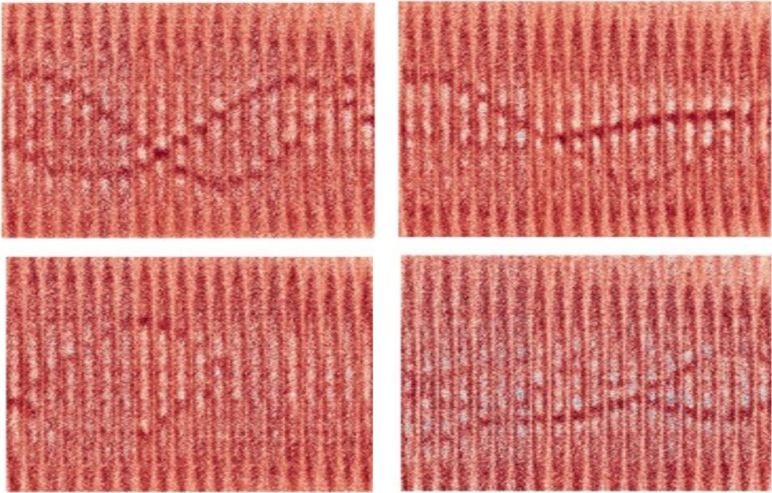
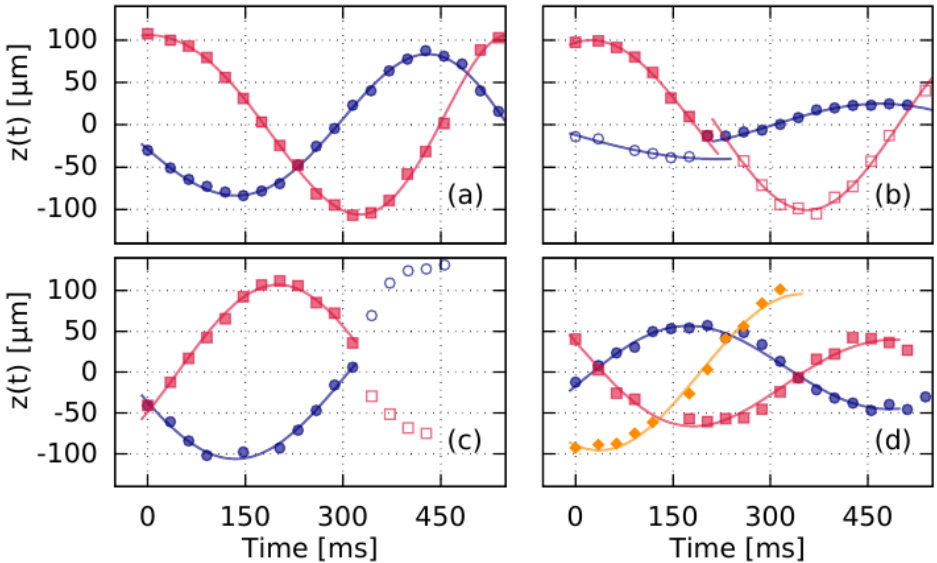
Interaction among vortices: lifetime measurement

1 or 2 vortices: decay by dissipation with the thermal fraction

3 vortices: faster decay



Interaction among vortices: phase delays

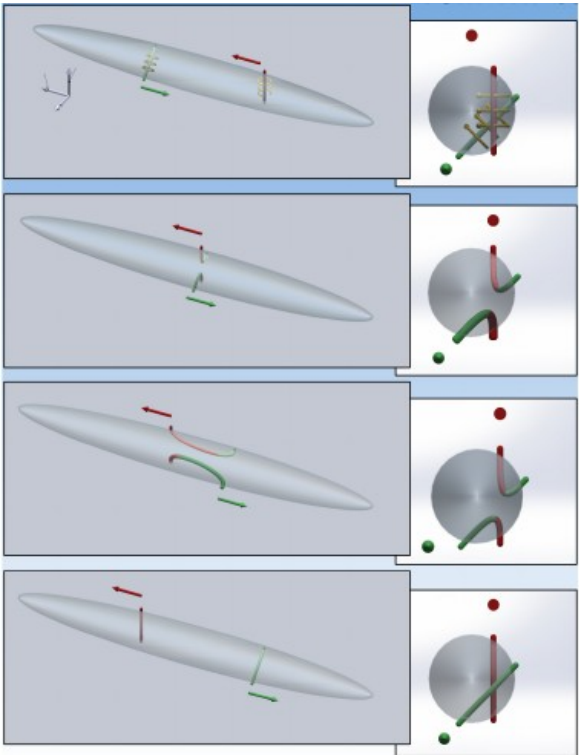


- Frequently: no visible interactions
- Frequently: change of visibility
- Sometimes: phase shifts
- Seldom: annihilations

Single reconnection energetically expensive due to nodal line stretching.

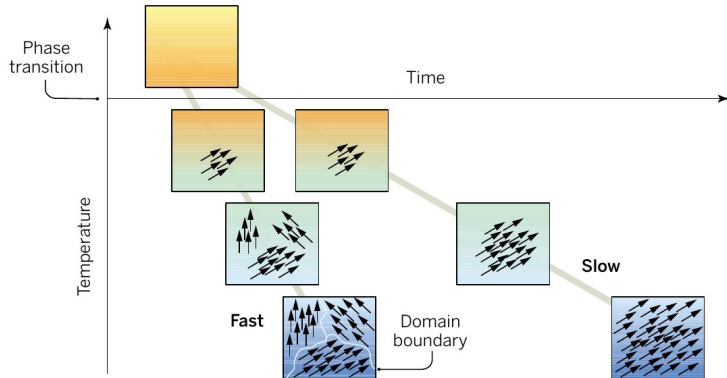
Possible alternatives:

- double reconnection
- rotation of the nodal lines when approaching

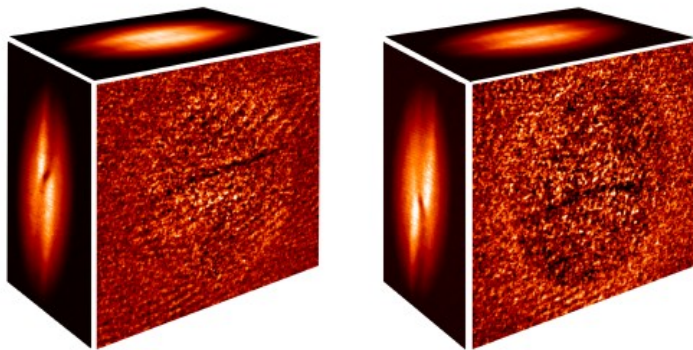


Summary

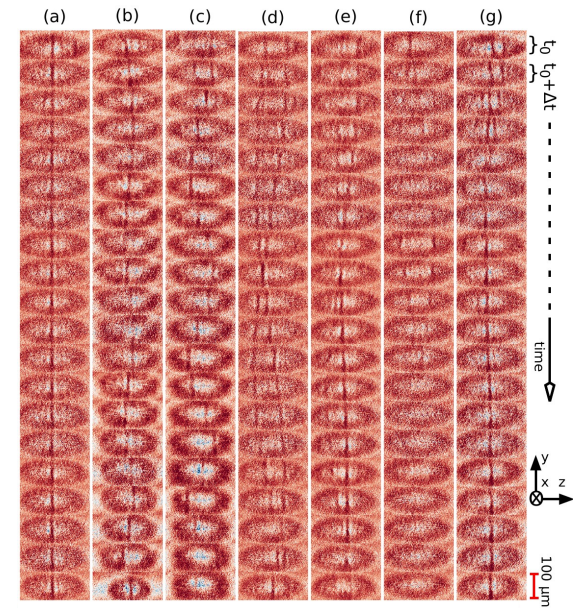
formation



nature



dynamics & interaction



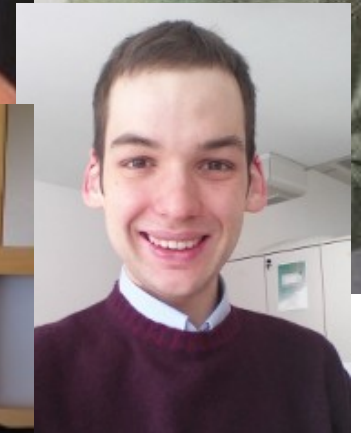
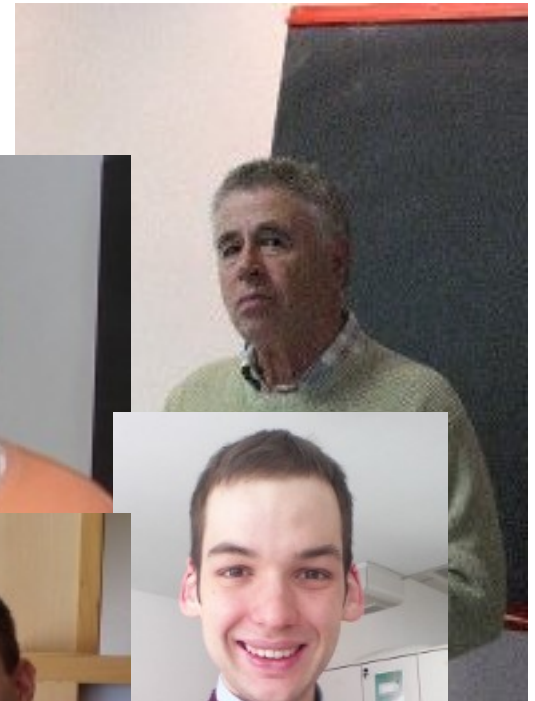
future developments

- investigation of post-quench dynamics after crossing phase transitions
- microscopic study of reconnection mechanisms

Thank you!



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